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Iowa State University of Science and Technology Ph.D., 1962 Chemistry, physical University Microfilms, Inc., Ann Arbor, Michigan

THEORETICAL ANALYSIS OF THE CHEMICAL BOND IN DIATOMIC HYDRIDE MOLECULES

by

Edgar Miller Layton, Jr.

A Dissertation Submitted to the Graduate Faculty in Partial Fulfillment of The Requirements for the Degree of DOCTOR OF PHILOSOPHY

Major Subject: Physical Chemistry

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INTRODUCTION

During the past decade there have appeared an increasing number of publications concerned with the rigorous calculation of molecular electronic energies. The impetus for this type of work has come from three main sources: the organization of formal self-consistent-field (SCF) programs for calculating molecular energies (4), the availability of general evaluation formulae for many types of molecular integrals, and the utilization of large computers for the systematic handling of the large number of required arithmetic operations.

As the mechanics of complicated molecular computations have become more and more standardized and automatic, a problem of interpretation has arisen. There are only a limited number of experimental quantities that can be compared with the theoretical calculations, such as molecular dipole moment, various ionization potentials and spectral transitions. While the agreement with experimental values is sometimes encouraging, it is particularly lacking as regards the binding energy. Furthermore, the binding energy represents only the totality of ell energy terms in the molecule, and

1Summary reviews are presented in References 1, 2 and 3.

 2 Typical integral projects are those of the Tokyo group (5) or the Chicago group (6, 7, 8 and 9).

hence provides little insight as to the detailed nature of chemical binding in the system. The population analysis of Mulliken (10), although in itself not based on the energy, is a noteworthy attempt to fill this gap. By this method, a qualitative estimate is made of relations between the distribution of electrons and binding effects.

Quite recently an energy analysis, based on density matrix properties rather than a particular type of wave function, has been developed by Ruedenberg (11) which is designed to gauge the quality of the various molecular calculations and to trace the energetics of bond formation. It attempts to extract specified physical and "geographical" energy quantities from different molecules in a prescribed It is conjectured that similarities will occur in manner. the more refined calculations for analogous atomic environments. Moreover, it is predicted that a mathematical construct termed interference will reflect bonding ability. The extended partitioning contained in this theory has been designed in the hope of accounting for the energy balance in molecule formation from isolated atoms, as well as bringing out those peculiarities which give the individuality to a particular molecule.

In the present investigation this method of analysis is applied to the linear-combination-of-atomic-orbitals molecular orbital (LCAO MO) calculations of a series of

hydrides, recently published by Ransil (12). The results indicate that it is possible to extract, from calculations of this type, a substantial amount of significant energetic information in addition to the total binding energy. This information provides a quantitative basis for the reasoning about contributory factors to chemical binding.

In spite of a broad variation in electronic parameters, trends can be established. The predicted importance of constructive interference, contractive promotion and sharing penetration is verified. Moreover, sets of valence atomic hybrid orbitals are found which are of consequence for the binding process.

While it is recognized that the wave functions analyzed are far from perfect, they offer the attractive possibility of performing the analysis on a completely analogous series of systems for which equivalent calculations have been executed.

Separate reports present the results of this analysis applied to the water molecule (13) (prototype analysis), the hydrogen molecule ion (14), and some diatomic homonuclear systems (15).

DESCRIPTION OF MOLECULAR SYSTEMS

The molecular systems to be considered are the diatomic hydrides of lithium. boron. nitrogen (for two values of the internuclear distance) and fluorine. The wave functions chosen are the SCF LCAO MO wave functions based on a minimal set of atomic orbitals, recently obtained by Ransil (12)? The MO's for each hydride are determined by the Hartree-Fock formalism (4) for three approximations concerning the orbital exponents. In the simplest case (denoted by SAO), the orbital exponents (ζ) are chosen by Slater's rules.⁴ In a more refined determination (BAO case), the **ζ'**s are found by a minimization of the separate atomic energies with respect to their **č**'s (17). In the third approach the **č**'s (BMAO) are specified by minimizing the molecular energy, at an experimental internuclear distance, with respect to all **ζ's.** The first approximation is the crudest, but has shown itself to be remarkably effective for many purposes. The second method satisfies the virial theorem separately for each of the combining atoms. The last type of calculation is the most satisfactory of the three and it would satisfy the virial

³For the execution of our analysis, values of all integrals, between all the atomic orbitals involved, are required. We are grateful to Dr. Ransil for providing us with this considerable amount of data.

⁴See, for example, p. 40 of Reference 16.

theorem for the molecule if the calculation were carried out at the theoretically predicted equilibrium distance. Unfortunately, this has not been done and the virial theorem is not satisfied even for this third approach. It is felt that the violation of the virial relationship for all cases presents a serious source of error for the promotion states and therefore precludes a completely satisfactory interpretation of the latter (11). The molecular energy obtained in the last method is, however, an improvement over both the SAO and BAO cases, and some distinct variations do appear in the analytical partitioning, as will be seen later.

Table 1 presents the internuclear distances and ς values for the hydride systems. The subscripts for the ς 's refer to the standard AO's: 1s on the hydrogen atom, hereafter referred to as h, and 1s, 2s, $2p\sigma$, $2p\pi$, and $2p\pi$ on the other atom, hereafter denoted by k, s, σ , π , and π respectively. These entries will be discussed in the text when appropriate.

The significant results for both NH bond lengths are contained in the calculations on the NH[•] system. The identifying supercript for the particular bond length chosen will be dropped henceforth and the shorter bond length will be implied throughout the analysis.

All wave functions represent ${}^{1}\Sigma$ electronic states of the molecule, i.e., a closed-shell configuration. This con-

System		R _{AB} (Angstrom units)	ζ _h	کر _{لا}	۲ ₅	ζσ	ζπ
LiH	SAO BAO BMAO	1.5953 ₅	1.0 1.0 .9766	2.7 2.6865 2.6909	•65 •6372 •7075	•65 •6372 •8449	
BH	SAO BAO BMAO	1.2325	1.0 1.0 1.1860	4•7 4•6794 4•6805	1.3 1.3383 1.2955	1.3 1.2106 1.3168	
NH *	SAO BAO BMAO	1.0455	1.0 1.0 1.4096	6.7 6.6652 6.6703	1.95 1.9236 1.9442	1.95 1.9170 2.0959	1.95 1.9170 1.7697
NH **	SAO BAO BMAO	1.1282	1.0 1.0 1.3574	6•7 6•6652 6•6707	1.95 1.9236 1.9307	1.95 1.9170 2.0236	1.95 1.9170 1.7 <i>5</i> 99
HF	SAO BAO BMAO	•9171	1.0 1.0 1.3163	8.7 8.6501 8.6533	2.6 2.5639 2.5551	2.6 2.5498 2.6693	2.6 2.5498 2.4965

Table 1. Molecular constants for the hydride systems

figuration corresponds to the experimental ground state of each system except NH. Here, the ground state is ${}^{3}\Sigma$, so that the calculations represent one of two known excited states. The combining nitrogen atom, in fact, possesses no lone-pair σ electrons. The eight molecular electrons are distributed in the following fashion: two in the nitrogen K-shell, two in a bonding hybrid composed of the nitrogen s, σ and hydrogen h orbitals and two in each of the two nitrogen π orbitals. This circumstance leads to certain deviations in the NH analysis with respect to other members of the series.

NATURE OF THE ANALYSIS

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A detailed description of the energy analysis is contained in the original article (11). A supplementary critical discussion of interpretative aspects is found in the report on the water molecule (13). The principal features of the analysis are briefly repeated here for the general diatomic system in order to facilitate an understanding of the details to be presented later.

The transition from isolated atoms to a molecule is considered purely from a conceptual viewpoint.⁵ The molecular electronic energy of a diatomic molecule (under the Born-Oppenheimer approximation, excluding magnetic and relativistic effects) may be written⁶

 $E = z_{A}z_{B}/R_{AB} + \int d\tau h \rho + \frac{1}{2} \int \int d\tau_{1} d\tau_{2} \pi / r_{12} , \quad (3.1)$

where A and B index the atoms, 1 and 2 index the electrons, ρ and π are the first- and second-order matrices, <u>h</u> is the

⁵The allied problem of the kinetic mechanism of molecular formation is another question. An interesting, but formidable, project would be the solution of the SCF equations for selected internuclear distances (ranging from infinite separation to an equilibrium position) and application of the present analysis to the resulting set of wave functions.

⁶Atomic units are employed. The conversion factor for the atomic unit of length is $1 a \cdot u = 0.5291$ A and for energy is $1 a \cdot u = 27.2052$ eV.

⁷See Reference 11 or 18 for a more complete treatment of density matrices and an expanded discussion of this equation.

one-electron operator including electronic kinetic energy and attractive nuclear-electronic potential energy and \underline{r}_{12} is the interelectronic distance. The density matrices are given as expansions in terms of atomic orbitals;

$$\rho = \sum_{\substack{r \in S}} \sum_{p(r|s)\chi_r(1)\chi_s(1)}, \qquad (3.2)$$

$$\pi = \sum_{r} \sum_{r} \sum_{s} \sum_{s} p(r\overline{r}|s\overline{s})\chi (1)\chi (1)\chi (2)\chi (2), \quad (3.3)$$

where r, \overline{r} , s and \overline{s} index the AO's (from either atom, in any sequence), and the p's are bond orders and pair-bond orders. The physical significance of the density matrices is established by the conservation relationships

$$\int d\boldsymbol{\tau} \boldsymbol{\rho} = N, \text{ the number of electrons }, \qquad (3.4)$$

$$\int d\boldsymbol{\tau}_1 d\boldsymbol{\tau}_2 \pi = \frac{1}{2} N(N-1), \text{ the number of}$$
(3.5) electron pairs.

The energy decomposition is effected by partitioning the density matrices, on the basis of considerations concerning physical interactions, and grouping these components with the corresponding integrals over the atomic orbitals. These energy fragments are then further apportioned to the appropriate atom, orbital or orbital pair along with a proportionate fraction of the nuclear-nuclear repulsion energy. All energy components remain separated in kinetic and potential contributions in order to trace the energetic balance occurring in the molecular formation process. At all steps, conservation restrictions analogous to 3.4 and 3.5 are applied and corresponding electronic populations are determined. To emphasize the molecular role of the atomic orbitals, an initial intra-atomic diagonalization of the bond order matrix is performed which quantitatively defines inner (K-shell), lone-pair and bonding hybrid orbitals. These are labelled the valence atomic orbitals (VAO's).

The electron density is partitioned into three components, viz.,

$$\boldsymbol{\rho} = \boldsymbol{\rho}^{\mathbf{CL}} + \boldsymbol{\rho}^{\mathbf{I}} = \boldsymbol{\rho}^{\mathbf{PR}} + \boldsymbol{\rho}^{\mathbf{T}} + \boldsymbol{\rho}^{\mathbf{I}}$$
(3.6)

where the superscript CL refers to the quasiclassical portion, I to the fraction participating in interference, PR to promotion and T to transfer. The interference is sufficiently novel and basic that further comment is warranted. This is that part of the electron density which distinguishes the quantum mechanical interactions in the process of bonding. Interference constitutes a spatial modification in the total electron density of the molecule without addition or deletion of charge density. The interference density between orbital a on atom A and orbital b on atom B is defined as

$$(Aa, Bb) = (Aa)(Bb) - \frac{1}{2}S_{Aa, Bb}[(Aa)^2 + (Bb)^2] .(3.7)$$

This function has the appearance of overlap dependence in ³ both terms. The interference effects are postulated as being crucially associated with bonding behavior.⁸

The electron pair-density partitioning is considerably more complex. Although it plays a substantial part in the ultimate analysis, its decomposition will only be intimated. Because of the Hartree-Fock SCF origin of the MO's, the pair density may be written

$$\pi(\underline{\mathbf{x}}_{1}\underline{\mathbf{x}}_{2}) = \rho(\underline{\mathbf{x}}_{1})\rho(\underline{\mathbf{x}}_{2}) - \frac{1}{2} \left|\rho(\underline{\mathbf{x}}_{1}|\underline{\mathbf{x}}_{2})\right|^{2} , \qquad (3.8)$$

or
$$\pi(\underline{x}_{1}\underline{x}_{2}) = \rho(\underline{x}_{1})\rho(\underline{x}_{2}) - \pi_{x}(\underline{x}_{1},\underline{x}_{2})$$
, (3.9)

where the first term represents a coulombic interaction and the second term exhibits exchange character. These equations give the explicit spatial dependence of having one member of the electron pair at point x_1 and the second member at x_2 . Upon substitution of 3.6 into 3.9, and isolation of the various π_x components, energy contributions from the second order density matrix may similarly be incorporated into the orbital and orbital pair decomposition noted above. More accurately, the combined coulombic and exchange fragments are rearranged into interference and interference-free parts. The interference-free parts, furthermore, contain the contri-

⁸See Reference 14 for a lucid explanation of the extremely fundamental nature of interference.

butions attributed to both the intra-atomic and interatomic sharing penetration categories. (See Appendix, especially Equation A.5, ff.)

Since the energy value of interest is that arising from the formation of a molecule, a logical extraction of energy relating to the isolated atomic ground states is desirable. This is easily accomplished so that the total molecular binding energy may be arranged in the form,

$$E_{MB} = E - E^{g} = \begin{bmatrix} E^{PRH} \\ E^{PRC} \end{bmatrix} + \begin{bmatrix} E^{QCN} + E^{SPN} + E^{SIN} \\ E^{QCT} + E^{SPT} + E^{SIT} \end{bmatrix}$$
(3.10)

where the superscripts PRH and PRC denote energy effects attributable to hybridization and contractive (clustering) promotion; QC, SP, and SI refer to energy terms from quasiclassical electrostatic, sharing penetration and sharing interference both before (N) and after (T) charge transfer. As demonstrated in Reference 13, the promotion and quasiclassical terms may be considered as non-sharing phenomena arising at infinite separation and at the equilibrium internuclear positions, respectively. The penetration and inter-

⁹The hybridization promotion is taken from the bond order density diagonalization, whereas clustering promotion is taken from the observed (14) increased electron density in the neighborhood of the nuclei. Alternatively, promotion may be described as angular (hybridization) and radial (contraction) density rearrangements.

ference terms may similarly be identified primarily with longrange and short-range sharing interactions.

All of the required arithmetic operations were executed on the Iowa State University Cyclone, an electrostatic memory digital computer patterned after the Illiac at the University of Illinois.

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GROSS CHARACTERISTICS OF BINDING ENERGY PARTITIONING

Qualitative Concepts

The approach followed here recognizes five major kinds of energetic interactions in the constitution of the molecular binding energy; promotional effects, quasiclassical effects, sharing penetration effects, sharing interference effects and charge transfer effects. These main divisions are again subjected to a variety of further subdivisions which will be examined in detail in the further course of the discussion. It appears, however, most convenient to first focus attention on the gross effects. For each molecular analysis in the Appendix, they are listed in the first table labelled "Binding Energy Partitioning". A comparative graphical representation of the main features for all molecules, in the SAO and BMAO approximations, is given in Figure 1. (The BAO approximation essentially duplicates the SAO behavior here and throughout the remainder of the analysis except where specific reference is made to the contrary.)

Three curves are given for each system. One represents energetic effects before "charge transfer", a second due to the characteristic charge transfer contributions and the third for the total of these. Each curve is displayed

Figure 1. Theoretical molecular binding energy decomposition, by intraand interatomic contributions, for SAO and BMAO approximations. PR stands for promotion, QC for quasiclassical, SP for sharing penetration and SI for sharing interference. Subdivisions indicate atomic or bond origin of contribution. Heavy atom indexes the first two promotion contributions, viz., hybridization (H) and contraction (C). Separate plots indicate total (\bigcirc), neutral (\bigcirc) and transfer (\triangle) categories.



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as a running total.¹⁰

The first three points represent promotion effects on the two atoms. Under the heavy atom symbol, two promotional mechanisms are distinguished: hybridization (denoted by H; hardly to be confused with the symbol for the hydrogen atom) and contraction (denoted by C; expansion being considered as negative contraction) which results from variations in the orbital exponents. The latter is the only kind of promotion for hydrogen and only applies to the BMAO approximation.

The next three points describe all those interactions which, somewhat casually, have been termed "quasiclassical" in the previous discussions of the present analysis (11, 13). In a more careful terminology, this name is a proper description only for the third value, which appears under the "bond" heading. This latter contains those interatomic potential energies which would result from a classical electrostatic energy calculation based on the electron density clouds before electron sharing. In contrast, the first two "quasiclassical" contributions arise entirely from charge transfer. They represent, for each atom, the intra-atomic energy changes resulting from the addition or loss of charge to the promotion state. For this reason, they will subsequently be denoted as intra-atomic charge-addition energies. The arbitrariness

¹⁰Note that the energy scale for NH is interrupted and more condensed, for reasons to be discussed below.

inherent in a definition of transfer, such as is used here, will further be discussed in later sections.

The subsequent three points represent the energetic contributions arising from sharing penetration. This section contains those electron repulsion effects which are specifically a consequence of the sharing of valence electrons between different atoms, as contrasted to those electron repulsions which are incorporated in the quasiclassical energy as shielding effects. These sharing penetration energies again have two intra-atomic portions and an interatomic portion.

The final category provides the interference energy, the most important of all, which is totally interatomic in character.

Quantitative Results

The justification of the present analysis lies in the fact that these enumerated categories exhibit characteristically similar behavior in each of the molecular systems analyzed. The only exception is the BMAO calculation of HF, where the separation of total effect into neutral and transfer parts does not follow the general pattern. A close examination indicates that a miscalculation must have occurred in the automatic calculation of this system, as will be

elaborated upon later, possibly connected with the inversion of some ill-conditioned matrices.

The "neutral plot" (i.e., contributions before transfer) always shows four typical features: a slight energy rise from hybridization and contraction promotion, a somewhat attractive¹¹ (1.0 to 3.0 eV) interatomic quasiclassical energy, a somewhat larger (typically 2 to 4 eV with NH having the extremes of 5 to 8 eV) bond-opposing rise due to sharing penetration and always the result of large intra-atomic effects dominating slightly smaller interatomic effects of opposite sign, and last, the significant energy drop furnished by interference. The last phenomenon must be considered as the source for covalent bonding, analogous to the traditional qualitative resonance stabilization in valence bond theory. It increases regularly with atomic number and so does the binding energy in the optimal BMAO case.

The transfer plot has no promotional contribution and a negligible contribution to interference. The overall effect is negative, indicating the advantage gained by the transfer of charge (excluding the SAO calculation of HF). As seen in

¹¹Here and in the following we denote a negative energy contribution as "attractive" and a positive energy contribution as "repulsive", i.e., attraction and repulsion are used synonymously with binding and anti-binding. This usage must not be confused with the attractive or repulsive quality of forces at the equilibrium or any other distance. Forces are not being discussed in the present investigation.

the next section, charge transfer is very small for the BMAO calculation of BH, relatively small in the SAO calculations of BH and HF, substantial for the SAO calculation of LiH and the BMAO calculations of LiH and HF, and very large for NH. This sequence is reflected in the transfer energies. In all cases, the electron addition energy is positive on the more electronegative partner of the molecule and negative on the other atom. Also, charge transfer invariably provides an attractive interatomic quasiclassical effect. The influence of transfer on the sharing penetration is generally opposite to that of the neutral effect. This is in agreement with the concept that charge transfer tends to suppress the degree of covalent electron sharing.

In all cases (except the disreputable BMAO approximation for HF), the major effects before transfer dominate the total values so that the plot of the totals is similar to, although not identical with, that of the neutral part.

Detailed physical explanations and interpretations for the qualitative behavior of the various aforementioned contributions have been given in the previous expositions of the present analysis (11, 13). The results reported in the preceding paragraphs are in complete agreement with the views and conjectures expressed in those discussions.

ELECTRONIC DISTRIBUTION

The second table of each molecular analysis, entitled "Density Partitioning for Valence AO's", gives information about the distribution of electrons in the individual molecules. The following conclusions can be drawn from the quantitative information contained in its sections.

Hybridization

The four molecules studied form two pairs: the hybrids of lithium and nitrogen show similarities as do those of boron and fluorine. The two sets differ significantly in the composition of the bonding and lone-pair s- σ hybrid VAO's which are formed for each heavy atom. The reason is that for boron and fluorine, the total lone-pair occupations are approximately two electrons while in lithium and nitrogen, the lone-pair orbitals are unoccupied. Since the lone-pair orbital is empty in lithium and nitrogen, the bonding orbital hybridizes purely for its own benefit. The s σ promotion energy, which must be weighed against binding effects, is entirely determined by the bonding orbital. Under these conditions the optimum binding orbital possesses mainly s character with a small but important admixture of σ character. Its orbital energy is, of course, much lower than that of the

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unused lone-pair orbital which primarily has σ character. In boron and fluorine, on the other hand, the promotion energy of the lone-pair orbital cannot be ignored, but must be kept small for two reasons: it contains two electrons, as compared with one in the bonding orbital, and moreover, it does not benefit from energy lowering by binding effects. For these reasons, the composition of the lone-pair and bonding orbitals are reversed; the lone-pair is predominantly s, whereas the bonding orbital has to be content mainly with σ character and a token s admixture.

Contraction and Expansion

This hybridization picture, in conjunction with the ideas of contraction promotion developed in References 11 and 14, explains the changes in orbital exponents of the valence orbitals as given in Table 2. In lithium and nitrogen, the s

Orbital	LiH	BH	NH	HF	
k	•0044	.0011	•0051	•0032	
S	.0703	0428	•0206	0088	
σ	.2077	.1062	•1789	•1195	
h	0234	•1860	•4096	•3163	

Table 2. $\Delta \subset$ values, BMAO compared to BAO

and σ orbitals, both being involved only in constructive interference via a bonding orbital, experience a contractive promotion. In boron and fluorine, the σ orbital, the major contributor to the bonding orbital, still contracts whereas the s orbital, primarily associated with the lone-pair orbital, and hence involved in destructive interference, shows expansion.

The hydrogen orbital in BH, where charge transfer is negligible, shows a change in the orbital exponent about identical with that in the H_2 molecule, as discussed in Reference 11, indicating that it is characterized by contractive promotion arising from binding interference. In the other systems, transfer of charge is another effect which influences the hydrogen orbital exponent. In lithium, with additional charge transferred to the hydrogen, electron repulsion causes a considerable expansion sufficient to not only annihilate the bonding contraction but even forcing an expansion, as compared to the free atom. With nitrogen and fluorine, charge transfer is reversed and hence the bonding contraction on hydrogen is reinforced. Nitrogen, with the larger charge transfer, shows the bigger effect.

The π orbitals consistently experience an expansive promotion as seen in Table 1. Presumably the presence of the attractive hydrogen nucleus offers the paired π electrons assistance in relieving their repulsion energy.

Transfer Populations, Bond Orders, and Valence Active Populations

Bonding and transfer properties of the orbital densities are displayed in Figure 2. Corresponding to the well known electronegativity behavior, charge transfer to the heavy atom increases from a negative value in lithium to positive values through fluorine, with that in boron effectively vanishing. The anomalously large nitrogen value is due to its previously mentioned peculiar promotion state where, because of the unoccupied low-energy lone-pair σ orbital, the bonding electrons are considerably less shielded from the nuclear attraction.

The valence active populations of the bonding orbitals follow the behavior of the bond order of these orbitals with the h orbital. These valence active populations are remarkably constant over the series while the bond orders increase slightly. The proper explanation of these trends is not simple. It is a fact that the interference energy, to be discussed below, markedly increases with atomic number. This is a consequence of the increasing steepness of the wave functions near the nuclei leading to an increasingly negative kinetic contribution to the interference energy regardless of overlap. This must be seen as the main reason for the increase in the experimental and theoretical (in the BMAO case) binding

Figure 2. Valence-orbital bond-orders (O), valence activities (\bigcirc), and charge transfer populations (\bigtriangleup) for SAO and BMAO approximations

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energy. However, because of the concomitant increase in non-bonded interference repulsions arising from σ and π lonepair electrons, the shortening of the bond lengths is still less than the shrinkage of the AO's caused by the increase in effective nuclear charge. For this reason, the overlap integral is found to decrease along the series. In spite of these considerable changes along the series, it seems as if there were a tendency to keep the valence active population approximately constant and that therefore the bond order, for the bonding orbital, shows a slight increase with atomic number.

The bond order for the lone-pair orbital is negative, corresponding to a non-bonded repulsion, and also increases in absolute value with atomic number. The conspicuously low value for the bonding orbital bond order in NH presumably reflects the fact that the effectiveness of interference (resonance stabilization) diminishes because there is a strong difference in electronegativities as is evidenced by the large charge transfer. This supposition is supported by the fact that the BMAO calculation, with considerably larger transfer, shows a further decrease in bond order.

It is intriguing that the bond orders for the innershell and lone-pair orbitals of <u>all</u> molecular systems, when plotted as a function of overlap, fall remarkably accurately on <u>one</u> curve which passes through the origin and is nearly

linear. These bond orders are negative (anti-binding character) and their overlap integrals have values less than 0.35. The bonding orbitals, on the other hand, with positive bond orders and overlap integrals greater than 0.44, show the different behavior implied above. Except for NH, a reciprocal dependence on overlap is seen; as the bond order decreases from 0.7 to 0.5, the overlap integral increases from 0.45 to 0.7.

PROMOTION

Basic Considerations

In order to find promotion energies, the ground state wave function must be assumed for each participating atom. These were chosen as follows.

The wave function for the hydrogen atom was taken as the 1s configuration with an effective charge of $\zeta = 1$.

For the other atoms, a single determinantal wave function corresponding to a ground state atomic configuration was selected. These were $1s^22s$ for lithium, $1s^22s^2p\sigma$ for boron, $1s^22s^22p\sigma 2p\pi 2p\pi$ for nitrogen and $1s^22s^22p\sigma 2p\pi^2 2p\pi^2$ for fluorine. In the case of the SAO calculations, the Slater orbital exponents were also used in the ground state wave function. For the BAO and BMAO calculations, the ground state wave functions were assumed to have the BAO orbital exponents.

Mixing of orbitals was thus the only promotion effect for the SAO and BAO calculations and this will be referred to as hybridization promotion. In the case of the BMAO calculation, in addition to hybridization promotion, there were the promotion effects from the changes in orbital exponents. These will be referred to as contractive or expansive promotion.

The choice of the ground state is really impractical for the nitrogen atom, since an excited state of the NH molecule is undergoing analysis. More appropriate would have been a state of the configuration $1s^{2}2s2\pi^{2}2\pi^{2}$. Before this could be done, the computer being used underwent major alterations. The energy differences between the promotion state and ground state in nitrogen are too large to be understood purely as "promotion energies". Only after subtraction of the suitable excitation energy could this interpretation be reasonable. The K-shell is of course unaffected by this excitation and its promotional effects correlate with those of the other atoms.

The promotion energies are given in the intra-atomic tables for the individual molecular analyses. The essential results for all molecules are collected, for comparison, in Figure 3. In the promotional energies, the individualities of the particular isolated atoms become a rather important factor and hence the comparison of different atoms is more complicated than for the characteristically molecular energy effects.

Hybridization

The first observation is that the hybridization promotion energies are remarkably similar in all three

Figure 3. Summary of orbital promotion energy decompositions into kinetic and potential parts for SAO approximation (hybridization only) and BMAO approximation (hybridization and contraction). Negative contributions indicated by solid overscore.

	LiH		BH		NH			HF				
	KIN	РОТ	тот	KIN	POT	тот	KIN	РОТ	тот	KIN	РОТ	тот
k	.76	, 98	. 21	1.01	1.94	. 93	1.54	2.59	1.05	. 53	. 80	. 27
SAO ^s	1.83	3.68	1.85	4.04	10.63	6.59	57.12	56.10	1.01	8.27	13.16	4.89
	1.95	3.06	1.11	2.41	6.86	4.45	49.92	52.22	2.30	10.68	13.47	2.79
п	1			ł			51.72	42.86	8.87	.00	. 08	. 08
TOTAL	. 88	. 36	.53	. 62	1.83	1.22	5.13	25.21	20.07	1.88	. 34	2.22
······								<u> </u>	<u></u> _			
k	. 94	1.22	. 27	.40	.81	. 41	1.65	2.86	1.20	. 79	1.37	.57
BMAO ^s	1.77	3.54	1.77	4.38	11.03	ს .6 5	54,20	30.37	23.83	8.74	26,21	17.49
HYBRID ^σ	1.78	2,85	1.07	3.18	8.00	4.82	48.29	51.29	3.01	11.87	24.28	12.41
π							49.99	31,21	18.78	.00	.03	. 03
TOTAL	. 96	. 54	.43	.80	2.22	1.42	4,16	22.09	17.94	2.33	3.26	5.59
	{]			ļ			1		
k	. 65	. 63	.01	. 28	. 24	. 04	1.89	1,58	. 31	1.58	1.35	.23
s	1.28	1.19	.10	4.57	4.20	. 38	1.77	1.60	.17	2.00	1.92	. 14
BMAO _σ	1.35	. 99	. 37	4.24	3.78	.45	.33	. 46	.13	9.62	9.58	.04
CONTR m	1						14.77	14.74	. 03	7.32	7.34	.02
TOTAL	3.28	2,81	.48	.06	.18	. 12	25.56	25.84	.28	5.50	5.67	.17
h	.63	. 64	.01	5,53	5.06	. 47	13.43	11.14	2,28	9.97	8.61	1.36
TOTAL	2.65	2.17	.49	5.47	4.88	. 59	12.13	14.70	2.56	4.47	2.94	1.53
				ĵ –			{			ſ		
ĸ	1.59	1.85	. 26	. 68	1.05	. 37	. 24	1.28	1.51	.79	. 02	. 80
S BMAO	.49	2.35	1.87	8.95	15.23	6.27	52,43	28.77	23.66	10.80	28.13	17.35
	3.13	5.84	. 10	1.42	11.78	4.37	47.96	50.83	2.88	21.49	33.86	12.37
	1		0.0		2 20		35.22	16,47	18,75	7.32	7.31	.01
TOTAL	4.24	5.54	.90	.80	2. 59	1.54	29.72	47.94	18.21	3.16	8.93	5.76
h	.63	. 04	.01	5.53	5.06	.47	$\frac{13.43}{17.30}$	11.14	2.28	9.97	8.61	1,36
TOTAL	3.01	2.70	.91	4.67	2.67	2.01	16.29	36,80	20,49	6.81	. 32	7.12

COMPARISON OF PROMOTION ENERGIES

approximations. For this reason, the BAO values are omitted. In general, removal of electronic charge from an orbital results in a decrease of kinetic energy and an increase in potential energy. The latter tends to be the larger effect (although not twice as large since we are concerned with shielded nuclear attraction for which the virial theorem does not hold) and hence the total energy contributions from that orbital is raised. Addition of charge to an orbital has the opposite effect.

The simplest example is the promotion (approximately 0.35 electron) from the s orbital to the s- σ hybrid orbital in lithium. Due to the nature of the orbitals, the energy decrease provided by the σ contribution does not compensate for the increase of the s contribution. A similar effect is observed in boron. The promotion charge is much smaller (ca. 0.1 electron) since the two lone-pair electrons still mainly occupy the s orbital. Nevertheless, the energy effect is larger than in lithium because of the larger effective nuclear charge of the orbital. Fluorine shows about the same promotion population as boron and the kinetic and potential energy contributions are correspondingly larger. The total orbital contribution to the promotion energy is smaller, however, since the difference between kinetic (proportional to ζ^2) and potential (proportional to ζ) contributions is less.
The values obtained for the nitrogen atom behave differently since here the change in electronic correlation becomes a dominant factor. In the ground state there are two paired s electrons and three unpaired π electrons while in the excited state there are an unpaired electron in the s- σ hybrid and four paired electrons in the π orbitals. Thus, the contribution of the s orbital decreases because of unpairing, the contribution from the σ orbital increases because of depopulation and the contribution from the π orbitals are positive because of pairing.

The non-negligible promotion contributions of the Kshell are due to two circumstances; the K-admixture to the valence hybrid orbitals and the deviation of the q^{N} 's from 2.000, the latter resulting from our definitions of charge transfer populations. In this context it is interesting to observe that boron and fluorine, where the lone-pair orbital is occupied and therefore has primarily s character, exhibit a not insubstantial k admixture to the lone-pair VAO.

Contraction and Expansion

The <u>total</u> promotional energies from orbital contraction and expansion are so small that the changes in orbital exponent might appear inconsequential to a superficial examination. The separate potential and kinetic

fragments indicate, however, that this would be a serious misjudgment. In all cases, the contraction or expansion gives rise to a substantial reapportioning of the kinetic and the potential energy parts.

The σ orbitals, primarily involved in bonding, always show an increase in kinetic and a slightly smaller decrease in potential energy. The same is shown by the s orbital wherever it is involved in bonding to a significant extent, specifically in lithium and nitrogen. Contrariwise, the s orbital shows a decrease in kinetic and an increase in potential in boron and fluorine where it is involved in anti-binding and, as mentioned earlier, shows an expansive promotion. The k orbital always contributes a not insubstantial decrease in potential and increase in kinetic energy similar to a binding orbital. Perhaps this arises because of the K-admixture to the bonding orbitals; more likely however, constructive interference favors a shift of intra-atomic energy from the potential to the kinetic category not only for those orbitals which produce constructive interference but also for others unless there exist strong opposing factors.

The hydrogen orbitals also show the characteristic shift of energy from potential to kinetic associated with the contractive promotion arising from constructive binding interference. The one exception is in LiH where charge transfer has forced the expansive promotion previously discussed.

The total LiH molecule significantly shows, however, a lowering of the potential and an increase of kinetic energy (considerably larger in absolute magnitude than the calculated binding energy) which must be assigned to an overall contractive promotion. The same is found in all other systems if the subtotal, excluding π orbital contributions, is These results demonstrate the intrinsic connecconsidered. tion between potential lowering, contractive promotion and covalent binding which has been repeatedly emphasized (11, 13, 14). There can be no question that the fulfillment of the virial theorem, if the calculations were to be carried out at the theoretical equilibrium internuclear distance, would also be brought about by the appropriate amount of contractive promotion.

For reasons discussed earlier, all π orbitals experience an expansive promotion. Correspondingly, their kinetic energy decreases and their potential energy increases. In fluorine, the effect of the π orbitals is insufficient to upset the total molecular balance established by the bonding σ framework. In nitrogen, the π orbitals do effect a reversal of the σ balance but the significance of this is hard to trace since we are dealing with an excited molecular state and moreover, with an ill-suited atomic reference state.

Total Promotion

It is interesting to observe that the sum of hybridization and contraction promotion (except for nitrogen with its anomalous hybridization and π expansion) always leads to a positive and regularly increasing kinetic contribution. The potential contribution shows the same negative value for lithium and boron, both of which are free of π electrons. The slightly positive fluorine potential value may be spurious because the potential energies of the hybridization promotion are presumably affected by the suspect behavior of sharing penetration, resulting from the aforementioned miscalculation, in the BMAO analysis for this molecule (see below).

INTRA-ATOMIC CHARGE ADDITION

The contributions discussed in this section are taken from the first column, termed QCT, in the tables of the intraatomic decompositions in the molecular analyses. In the present method, the charge addition energies displayed in Figure 4 play a role comparable to that played in semiempirical thinking by the ionization potentials and electron affinity of the promoted state. The numerical values are very different, however, for two reasons. First, the addition energies used here are those calculated for the transferred charge with a one-electron hamiltonian containing an effective potential of the neutral promoted atom. Second. it does not contain the intra-atomic energy changes which charge transfer causes by modifying the sharing situation. As a result, the addition energies always are positive if charge is added and negative if charge is lost.

Within this context the charge addition energies show great regularity. In all cases the bonding orbital contributes the major effect. It closely parallels the charge transfer populations shown in Figure 2. The addition energy is negative for lithium, vanishing for boron, positive for fluorine and exceptionally large and positive for nitrogen. Invariably there is a substantial cancellation between kinetic and potential parts with the kinetic part determining the

Figure 4. Intra-atomic valence-orbital charge-addition energy contributions for SAO and BMAO approximations. \bigcirc , total; \bigcirc , kinetic; \triangle , potential

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final value. The effect on the hydrogen is generally opposite to that on the heavy atom. The enormous values for the kinetic and potential energy parts for nitrogen (compared with those for the other atoms) clearly exhibit the unusual character of this molecule in the present series in that here the heavy atom has a low-lying empty orbital. Since there is a very large cancellation between kinetic and potential energy in nitrogen, it is remarkable that the values of the totals still closely reflect the transfer populations.

QUASICLASSICAL INTERATOMIC ENERGIES

The interatomic quasiclassical electrostatic energies between orbital pairs are given in the first three columns of the table of interatomic contributions in the molecular analyses. The significant values are plotted in Figure 5. The total interaction "before transfer" is broken down according to interactions between the neutral hydrogen atom and fragments of the neutral heavy atom which are associated with its valence orbitals. Such a neutral orbital fragment consists of the promotion density of that orbital plus an exactly counterbalancing fraction of the positive nucleus. The contribution of one orbital pair to the total transfer energy is the sum of three interactions: the interaction between the neutral orbital fragment of one partner with the transfer density of the other partner, the corresponding interaction between the neutral orbital fragment of the second partner with the transfer density of the first, and the interaction between the transfer densities of the orbitals.

bh Orbital Pair

The interactions between the binding heavy-atom orbital and the hydrogen orbital represents the dominant molecular effect. Consider first the interaction between the

Figure 5. Interatomic quasiclassical energy contributions, and decomposition by valence-orbital pairs, for SAO and BMAO approximations. \bigcirc , total; \bigcirc , neutral; \triangle , transfer.

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neutral hydrogen atom and that part of the neutral heavy atom which we ascribe to the bonding orbital. In all cases, this neutral bonding-orbital fragment is attracted by the neutral hydrogen atom, because the latter attracts the closer, though diffuse, electronic charge of the bonding orbital more strongly than it repels the more distant nuclear fraction associated with the bonding orbital to form the neutral fragment. It is not surprising that the σ type bonding orbitals of boron and fluorine, approaching closer to the hydrogen atom, lead to a stronger shielded nuclear attraction than the s type orbitals of lithium and nitrogen. Thus. hybridization differences explain the saw-tooth appearance of the neutral bh interaction in Figure 5. The slight overall negative slope in the SAO case is presumably caused by the decrease in internuclear distance, although this effect is certainly moderated by the increasing shielding of the orbitals because of their shrinking with atomic number. In the BMAO case, this moderation is accentuated by the contractive promotion previously discussed so that, aside from the hybridization variations, the neutral interatomic interactions essentially remain constant over the series.

Except for nitrogen, the <u>transfer</u> effects of the bh pair are smaller than the neutral effects just discussed. This is so because of the small amounts of charge involved in the transfer process. The quasiclassical transfer energies

are negative except for boron. Their variations follow very closely the absolute values of the transfer populations, because the dominant contribution is the attraction between the transfer populations of the two orbitals. The deviation of boron arises because, here, charge transfer is so small that the other two contributions become important, namely the interaction between the transfer population in the heavyatom bonding orbital with the neutral hydrogen atom and the interaction of the transfer population on the hydrogen atom with the neutral bonding orbital fragment of the heavy atom. Both are repulsive in lithium and boron and attractive in nitrogen and fluorine; the former because the neutral hydrogen still attracts negative charge, the latter because the interaction of the hydrogen transfer population with the charge in the bonding orbital is dominant over that with the more distant nuclear fraction (see the similar discussion for the neutral quasiclassical energies). The latter contribution is, of course, the smaller effect since it arises from two diffuse charges. In boron, these two positive contributions overcome the attraction between the diminutive transfer populations.

For the reasons stated, the variations in transfer interactions follow a characteristically different pattern than those of the neutral energies. The totals show a rather smooth energy decrease along the series except for the

prominent deviation of the abnormal nitrogen charge transfer.

1h Orbital Pair

Consider now the lone-pair of the heavy atom. The interaction between the neutral lone-pair fragment and the neutral hydrogen atom is positive because the lone-pair orbital cannot sufficiently shield the repulsion between its associated nuclear fraction and the hydrogen atom. This repulsion increases with atomic number for the SAO case and is constant for the BMAO case, that is, it behaves like the absolute value of the quasiclassical neutral attraction found for the bh orbital pair, and indeed, for the same reasons.

In contrast to the situation for the bonding orbital, the transfer portions follow the transfer populations of the heavy atom, with their given signs. The explanation is that the major contribution is now the interaction between the transfer population on hydrogen and the neutral lone-pair fragment of the heavy atom which, because of insufficient shielding, is in turn dominated by the interaction with the positive nuclear fraction of the heavy atom. Thus boron, where electrons move to the hydrogen, exhibits attraction and fluorine, with a positive hole on the hydrogen, displays repulsion. The other two contributions to transfer energy involve, on the heavy atom, only the transfer population of

the lone-pair orbital. Because the transfer populations of the lone-pair are rather small and because the lone-pair orbital is more distant from the hydrogen atom, these contributions are considerably smaller than the one discussed first which involves the heavy atom nucleus.

nh Orbital Pair

For the π orbitals in nitrogen and fluorine, the transfer energies consist only of the interaction between the hydrogen transfer population with the neutral π orbital fragment of the heavy atom. Again, the repulsion of the positive nuclear part of the fragment is insufficiently counterbalanced by the attraction of the π electrons giving positive transfer energy, obviously larger in nitrogen than in fluorine. In fluorine the π quasiclassical energies are lower than those of the σ lone-pair since shielding of the latter is clearly less than that of the former.

Total

Thus, it is evident that the repulsive quasiclassical energy contributions of the σ and π lone-pair orbitals temper the attractive contributions of the bonding orbital. The total molecular values for the quasiclassical energies still

show the principal features of the bh interaction but to a more moderate degree. The neutral values, while remaining remarkably constant over the series, reflect the effect of hybridization. For the transfer values, the influence of the lone-pair interactions is more consequential and in fluorine even leads to a transfer repulsion. It is noteworthy that in both calculations, cancellation of effects from various orbital pairs leaves the total molecular values all negative and confined to a remarkably narrow range.

SHARING PENETRATION ENERGIES

The energy values arising from sharing penetration are given in the tables of intra-atomic and interatomic contributions for the molecular analyses. Since penetration is a single phenomenon, both intra- and interatomic contributions must be considered simultaneously and hence the sum of both is given for each orbital pair in Figure 6^{12} In discussing this figure, one must remember that the construction of sharing penetration energies is the most questionable part of the whole analysis. It is therefore gratifying that the results show a reasonable consistency.

In all cases, except one, the sharing penetration energy is positive in complete agreement with the general ideas of the analysis, namely that sharing increases the penetration of electrons from different atoms and hence leads to a net increase of electronic repulsion. The exception is lithium where the sharing penetration energy is slightly negative for which there is no obvious explanation at present. Note that these contributions are very small, however. In

¹²In arriving at these figures, it was necessary to partition the intra-atomic hydrogen contribution according to the heavy atom orbital. For the neutral sharing penetration, this pro-rating was done by use of the interatomic values of the exchange pair-population contributions, which table is given. For the transfer sharing penetration, the separation was carried out according to the orbital charge transfer populations.

Figure 6. Molecular sharing penetration energy contributions and decomposition by valence-orbital pairs for SAO and BMAO approximations. \bigcirc , total; \bigcirc , neutral; \triangle , transfer.



the BMAO case of fluorine, which shows a <u>large</u> negative sharing penetration energy, there appears to have occurred a mechanical computational error. This is indicated, in the table showing the exchange contributions to the pair populations, by the fact that the total interatomic value of the lhorbital pair (-0.062) is divided into unreasonable neutral (0.193) and transfer (-0.254) portions. Most likely the matrix inversion required for the transfer penetration effects was disturbed by ill-conditioning.

The lone-pair orbital, with the smaller binding interactions, shows the smallest penetration energies, by far. For the bh orbital pair, the total sharing penetration energy increases very regularly from lithium to fluorine. This increase is entirely attributable to the shrinkage of the heavy atom orbitals whose intra-atomic contributions are the dominant terms. This regularity is not entirely a trivial consequence of the manner in which these numbers were obtained since the total values were calculated by adding together the parts due to transfer and those from neutral penetration, both of which are computed separately. The transfer contributions essentially follow the absolute value of the transfer populations; the neutral contributions seem to reflect the hybridization modulation encountered repeatedly, although it is difficult to trace a theoretical connection between the two.

SHARING INTERFERENCE ENERGIES

The interference energies, corresponding to the resonance stabilization operative in the molecule, are also found in the table of interatomic energy contributions in the molecular analyses. The salient features of the "before transfer" interactions are plotted in Figure 7. The modifications from charge transfer are too small to be of consequence. In complete agreement with the basic point of view underlying the present investigation, it is found that the binding bh orbital pair always shows a negative interference energy arising from a negative kinetic part and in spite of a positive potential part, both features being characteristic of constructive interference. The converse holds for the anti-binding Lh orbital pair which is characterized by destructive interference.

bh Orbital Pair

The kinetic energy part of the bh pair decreases regularly and rapidly to large negative values. Since we have seen that the bond orders change very little, it is clear that this drop originates in the kinetic resonance integral. In fact, it is a consequence of the increase in gradient of the bonding orbital near the nucleus as the

Figure 7. Molecular sharing interference (neutral) energy and contributions by orbital-pairs for SAO and BMAO approximations. \bullet , total; \bigcirc , kinetic; \triangle , potential.

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atomic number increases. (The interference energy is proportional to the square of the gradient which, in turn, is proportional to the effective charge of the orbital.) This increase in steepness is not only derived from the increase in effective charge of the s and σ orbitals, but also from the even more drastic increase in effective charge of the orthogonalizing k admixture to the L-shell orbitals. (The variation in overlap integral with atomic number is much less consequential.) On the other hand, the increase of potential interference with atomic number arises from the increase of actual nuclear charge which renders the shift of charge, from the atomic region into the bond region, more costly.

The potential part is smaller than the kinetic part and apparently more susceptible to hybridization effects. One would expect the bonding hybrids of lithium and nitrogen to show stronger relative interference effects because their hybridization is entirely determined by the interactions of the bonding orbital without disturbance from the anti-binding σ lone-pair orbital. In the SAO case, we do indeed find that the boron and fluorine kinetic values lie slightly above the line joining the lithium and nitrogen values. Correspondingly, the line joining the potential values of boron and fluorine is below that joining lithium and nitrogen. (Thus, the equality in lithium and boron potential bh interference

energies, for the SAO case, is explained by the fact that boron has less favorable bonding hybridization but a stronger effective nuclear charge.)

With the exception of nitrogen, the interference effect is accentuated in the BMAO case. This effect of contractive promotion is comparable to the increase in interference energy with atomic number. The deviation of nitrogen is accidental since the somewhat arbitrary SAO ζ sufficiently varies from the BAO ξ (for the s orbital) so that there is actually a slight orbital expansion, in going from the SAO to the BMAO situations. If the BAO values are compared with the BMAO values, the behavior of nitrogen falls in line with that of the other atoms although the not inconsiderable decrease in bond order (see Figure 2) markedly diminishes the effect arising from the orbital pair still reflects the difference in hybridization.

lh Orbital Pair

The interference energy of the *l*h pair is anti-binding due to its positive kinetic part which shows the expected increase from boron to fluorine as a consequence of orbital shrinkage analogous to the decrease of the kinetic part in the bh pair. Similarly, the expansive promotion of this

lone-pair orbital, referred to earlier, is the reason for the BMAO values being smaller than their SAO counterparts. Ås usual. the potential contributions of destructive interference are negative but they also represent a much smaller fraction of the total interference energy in the 1 h pair as compared to the bh pair. Such behavior seems to be characteristic of the interaction of a lone-pair orbital on one atom and a bonding orbital on another. As has been discussed in detail in the case of the water molecule (13), where a similar observation has been made, the explanation is presumably that here the destructive interference does not shift charge nearer to the nucleus but rather from one side of the nucleus (the bonding side) to the other side (the lone-pair side) 13 In fact, the potential energy is negligibly small in the case of boron.

πh Orbital Pair

The interference contributions for the π h orbital pairs arise solely from the exchange part of the second order density. As in other cases (13, 15), they are found to be small and negative and are suspected of being related to the

¹³In the present case this can be directly attributable to the unusually large k admixture to the boron lone-pair orbital.

energy lowering operative in Hund's rules.

Total

Similar to the situation in the case of the quasiclassical interactions, the total molecular effect is a moderated image of the bh interactions. The significant trend is the accentuation of total interference with atomic number upon which there is superimposed modulation which reflects hybridization.

CONCLUSIONS

Following the detailed discussion of the particular molecules and their behavior in the foregoing analysis, it is now appropriate to take a general overview of the results of the investigation.

A first question is what can be said, in view of the detailed analysis, as a general characterization of the electron pair bond in the hydrides considered. The differences between the four molecules have been seen to arise from two properties; the occupation of the σ lone-pair orbital and the electronegativity of the heavy atom.

The σ lone-pair occupancy (zero in LiH and NH, two in BH and HF) determines the hybridization of the lone-pair and bonding orbitals. This in turn is related to the characteristic variations in the neutral quasiclassical interactions favoring the situation with no σ lone-pair electrons. The increase in electronegativity with increasing nuclear charge is of course responsible for the variations in charge transfer along the series, i.e., (Li⁺H⁻), (B^OH^O), (N⁻H⁺), (F⁻H⁺). Associated with this is a corresponding variation of the charge transfer contribution to the binding energy. This effect is abnormally large in NH because of the empty σ lone-pair orbital.

A second effect of the increasing electronegativity

is, however, the regular increase in the interference energy via the kinetic energy of the progressively contracting orbitals. This is the main reason for the increase in binding energy along the series. It must be noted, however, that the anti-binding interference interaction between the lone-pair orbitals and the hydrogen atom rises similarly, and for this reason the increase in binding energy is not accompanied by a corresponding decrease in bond length. More precisely: The <u>actual</u> decrease in bond lengths along the series is slightly less than the shrinkage in orbital diameter due to increasing electronegativity.

As a consequence, the overlap for the various orbital pairs generally remains in the same range, with a slight decrease as atomic charge increases. This must be the reason for the bond orders, and more specifically, the valence active populations, remaining approximately constant and furthermore for the neutral quasiclassical interactions also remaining about constant, aside from the aforementioned hybridization modulation.

In spite of the differences in hybridization, the only trend discernible in the hybridization promotion energies appears to be the steady increase due to increasing the effective charge. The contractive promotion, while small in total, always embodies a substantial reapportionment of the energy from the potential to the kinetic type.

A second question is whether the proposed energetic interactions and the general technique of analysis appear to be justified for these hydride systems. Promotion of isolated atoms has been found expensive to varying degrees. (This is not a trivial observation. The choice of an appropriate reference state is of definite importance, but even a less desirable reference state indicates significant, and justifiable, internal alteration depending upon the molecule which eventually is to be formed.) In all cases the importance of orbital contraction to bonding is evident. Quasiclassical electrostatic interactions are usually of a slightly binding nature. There is an extreme cancellation between repulsive intra-atomic and attractive interatomic sharing penetration effects with a small net repulsion. The predicted importance of constructive interference, kinetic in origin, in forming a bonding situation is emphatically confirmed. Charge transfer effects are seen to improve the binding situation. In all cases, moreover, the observed characteristics and deviations could be understood with the use of only a few variables, such as hybridization, contraction and effective nuclear charge. A pleasant dividend has been the formulation of valence hybrid atomic orbitals which seem to play a fundamental role, not only in determining critical electronic populations, but also in defining the relative contributions to the different energy effects. The

obvious failings which do appear are generally localized (except for the interrelationship mentioned between promotion and sharing penetration) and interpretable. Under the present critical testing, the proposed energetic characteristics of molecular formation have generally been vindicated.

It is, of course, evident that further refinements of the present analysis would lead to a better description of interesting details. For such a closer examination of the method, the wave functions used here are presumably inadequate. Minimization with respect to orbital exponents, at the theoretically deduced equilibrium positions, is indispensable for a proper reapportionment of the kinetic and potential energy. The use of an extended basis to approach the true SCF wave function would also be desirable. The execution of the analysis for such a near-SCF solution as a function of internuclear distance would be worthwhile, before proceeding to more complex wavefunctions incorporating correlation features.

The implied inadequacies in the wave functions and the method of analysis should be kept in mind when judging the reported conclusions.

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APPENDIX: PRESENTATION OF NUMERICAL RESULTS

General Description

Each of the figures, 8 through 22, represents the complete molecular analysis for one of the fifteen hydride calculations. The molecules are arranged in order of increasing molecular weight. For each molecule, the approximations are listed in the order: SAO, BAO, BMAO. The first three parts are the summary binding energy totals; VAO decompositions, bond orders and populations; and the exchange contributions to the pair populations. The last two parts are the detailed decomposition of the energy fragments, by orbital and orbital pair, for each energy mechanism.

Summary Properties

The binding energy partitioning summary contains totals by atom, bond and molecule for the various energy mechanisms. The headings follow the notation of Eq'n 3.10. The entries are sums of the quantities discussed below for the explicit orbital and orbital pair decomposition.

The table of density partitioning shows the AO decomposition, the overlap integrals with the hydrogen ls
orbital and the populations and bond orders of the orthonormal VAO's. The q's are the total electronic orbital populations in terms of the valence inactive (p) and valence active (v) fragments, and arising from promotion (N) and charge transfer (T) sources. For review purposes, the following relationships are listed for the diatomic hydride VAO's:

$$q(r) = p(r) + v(r), \quad q(r) = q^{N}(r) + q^{T}(r);$$
 (A.1)

$$\Sigma q(r) = N$$
 (sum over all occupied orbitals); (A.2)
r

$$\sum_{r \neq h} q^{T}(r) = -q^{T}(h), \quad \sum_{r \neq h} v(r) = v(h); \quad (A.3)$$

$$v(r) = p_{(r,h)}S_{(r,h)}$$
 (A.4)

The table of exchange contributions to the pair populations lists coefficients, $q_x(Aa,Bb)$, which are used in conjunction with Eq'n 3.9. When a particular AO basis has been chosen (and consequently a particular VAO basis), the pair population between orbitals Aa and Bb is (as is shown in Reference 11)

$$q(Aa,Bb) = q(Aa)q(Bb) - q_x(Aa,Bb).$$
 (A.5)

Pair populations, as defined in A.5, are an interference-free property of the wave function. In the molecular energy, they

are the coefficients for a limited set of the two-electron integrals, i.e., those particular one- or two-center coulomb integrals having one electron totally in one crbital and the other electron totally in a second orbital, viz.,

$$[Aa2|Bb2]$$
, where A may equal B and a may equal b.

The entries in this table are the numerical values of q_X , corresponding to the interference-free parts deriving from promotion (PR), sharing penetration before (SPN) and after (SPT) charge transfer, and their total (TOT). That is,

$$q_{\mathbf{x}}(Aa, Bb) = q_{\mathbf{x}}^{TOT}(Aa, Bb)$$

1

=
$$q_X^{PR}(Aa,Bb) + q_X^{SPN}(Aa,Bb) + q_X^{SPT}(Aa,Bb)$$
. (A.6)

Again, the following particular relations are listed for these quantities:

$$\sum_{\mathbf{r}} q_{\mathbf{x}}^{\mathrm{PR}}(\mathbf{r}, \overline{\mathbf{r}}) = q^{\mathrm{N}}(\overline{\mathbf{r}}) \qquad (A.7)$$

$$\sum_{r \neq h} q_x^{\text{SPN}}(r, \overline{r} \neq h) = - q_x^{\text{SPN}}(\overline{r}, h)$$
(A.8)

$$\sum_{\mathbf{r}} q_{\mathbf{x}}^{\text{SPT}}(\mathbf{r}, \overline{\mathbf{r}}) = q^{\text{T}}(\overline{\mathbf{r}}) \qquad (A.9)$$

$$\sum_{\mathbf{r}} q_{\mathbf{x}}^{\text{TOT}}(\mathbf{r}, \overline{\mathbf{r}}) = q(\overline{\mathbf{r}}) \qquad (A.10)$$

$$\sum_{\mathbf{r}} \sum_{\mathbf{r}} q_{\mathbf{x}}^{\text{PR}}(\mathbf{r}, \overline{\mathbf{r}}) = N \quad (\text{sum over all occupied orbitals}) \quad (A.11)$$

$$\sum_{\mathbf{r}} \sum_{\mathbf{r}} q_{\mathbf{x}}^{\text{SPN}}(\mathbf{r}, \overline{\mathbf{r}}) = 0$$

$$\sum_{\mathbf{r}} \sum_{\mathbf{r}} q_{\mathbf{x}}^{\text{SPT}}(\mathbf{r}, \overline{\mathbf{r}}) = 0 \qquad (A.13)$$

$$\sum_{\mathbf{r}} \sum_{\mathbf{r}} q_{\mathbf{x}}^{\text{TOT}}(\mathbf{r}, \overline{\mathbf{r}}) = N \qquad (A.14)$$

Orbital and Orbital Pair Decompositions

The last two sections of each particular figure give the separate intra-atomic and interatomic energy contributions to the molecular binding energy. Under intra-atomic contributions there is an additional division between promotion effects arising in the original AO's and interaction effects occurring in the VAO's. Each of these contributions is further subdivided into kinetic (KIN) and potential (SNA for shielded nuclear attraction, NA for nuclear attraction and OEI for other electronic interactions) contributions.

Formulae for the subdivision of promotion energy into

orbital contributions are developed in detail in Reference 15. The criterion for this orbital partitioning is the difference in orbital electronic population for the promoted atomic valence state with respect to the atomic ground state. The promotion energy terms distributed to the orbitals are obtained from the intra-atomic integrals and coefficients of the density matrices. The total promotion energies are given by Equations A.15 - A.18:

SAO:
$$E_{HP} = E^P - E^g$$

$$= (p^{P} - p^{g}) \cdot \underline{1}(SAO) + (\pi^{P} - \pi^{g}) \cdot \underline{2}(SAO)$$
 (A.15)

BAO:
$$E_{HP} = (p^{P} - pg) \cdot \underline{1}(BAO) + (\pi^{P} - \pi g) \cdot \underline{2}(BAO)$$
 (A.16)

BMAO:
$$E_{HP} = (p^P - p^g) \cdot \underline{1}(BAO) + (\pi^P - \pi^g) \cdot \underline{2}(BAO)$$
 (A.17)

$$E_{CP} = p^{P} \cdot [\underline{1}(BMAO) - \underline{1}(BAO)] + \pi^{P} \cdot [\underline{2}(BMAO) - \underline{2}(BAO)]$$
(A.18)

In these equations, superscripts P and g refer to the promoted and ground states, <u>1</u> and <u>2</u> denote the one- and twoelectron intra-atomic integrals obtained by the specified approximation and HP and CP label hybridization and clustering promotion. Note that clustering promotion is only identified with the BMAO case. This results from the consideration that only in the BMAO approximation is there an energy minimization with respect to the molecular values of the ζ 's. By definition there is non-zero promotion in the hydrogen atoms only in the sense of contraction.

The remaining intra-atomic contributions, characteristic of the atom in the molecule, are reported for the VAO's. The QCT entries apportion the quasiclassical energy of the transferred charge to the appropriate orbital. Sharing penetration, both neutral and transfer, is separated in a similar fashion, with the fragments appearing only in the OEI rows, re-emphasizing the interpretation (13) that this effect arises from rearrangements within the electronpair density matrix.

All interatomic contributions are reported for the valence hybrid orbital pairs. In the QCT column, the SNA row is further separated into (left entry) the energy of the charge transferred to the heavy atom orbital in the field of the neutral shielded hydrogen nucleus and (right entry) the energy of the charge transferred to hydrogen in the field ascribed to the neutral heavy atom orbital fragment. The SIN column is similarly decomposed, the left entry showing the energy of the orbital-pair interference density in the field of the shielded heavy atom nucleus and the right entry giving the orbital-pair interference density energy in the field of

the shielded hydrogen nucleus. The SIT column gives the interaction energy of the particular interference density with the charge transferred to the particular heavy atom orbital (left) and with the charge transferred to the hydrogen atom (right). The SPN and SPT entries again illustrate the pair-density origin of this interaction.

The initial summary table may now be recognized as the totals at the bottom of each column of the intra-atomic (atom) and interatomic (bond) contributions with combined SNA entries.

Summarizing Explanation for Figures 8 - 22

The first section, labelled "Binding Energy Partitioning", lists the binding energy fragments by atoms and bond, decomposed in terms of kinetic interactions (KIN) and potential interactions. For the intra-atomic promotion contributions, the latter are subdivided into nuclear attraction (NA) and other electronic interaction (OEI). For all other intra- and interatomic categories the potential contributions are decomposed into shielded nuclear attraction (SNA) and other electronic interaction (OEI). Note that nonzero contraction promotion only occurs for the BMAO cases.

The second section, labelled "Density Partitioning for Valence AO's", gives VAO decompositions in terms of

basis functions; the VAO overlap with the hydrogen is orbital; valence inactive (p), valence active (v), total (q), promoted (q^N), and charge transfer (q^T) populations for these VAO's; and their bond order with the hydrogen 1s orbital.

The third section, entitled "Exchange Contributions to Pair Populations", gives the interference-free intra- and interatomic orbital pair coefficients of $\pi_{X}(\underline{x}_{1},\underline{x}_{2})$ and their decompositions in terms of contributions due to promotion (PR), sharing penetration before (SPN) and following (SPT) charge transfer.

The last two tables give the detailed partitioning according to orbitals and orbital pairs of the intra- and interatomic summary presented in the first table ("Binding Energy Partitioning") given above. In the SNA rows of the interatomic QCT column, the left entry represents the interaction of the neutral hydrogen with the transfer population of the heavy atom orbital, whereas the right entry represents the interaction of the neutral heavy-atom orbital-fragment with the transfer population on hydrogen. In the SNA rows of the interatomic SIN and SIT columns, the left entry gives the potential interaction of the interference density with the heavy atom while the right entry gives that interaction with the hydrogen atom.

		Promotion		Quasi- classical		Shan Penet	Sharing Penetration		ring	
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	
	KIN	. 85	-		-2.17					-1.29
• •	NA	-1.05	-		1.34					. 29
L1	OEI	. 69	-		. 37	4.65	-1.68			4.03
	TOT	. 53	-		46	4.65	-1.68			3.04
	KIN	_	-		4.21					4, 21
••	SNA	-	-		-3.16					- 3, 16
н	OEI	-	-		. 82	10.15	-8.98			1,99
	TOT	-	-		1.87	10.15	-8.98			3.04
	KIN							- 3. 0 3		- 3. 0 3
DOND	SNA			99	. 47			1.32	.00	. 80
BOND	OEI				91	-11.34	7.09	03		-5.19
	тот			99	44	-11.34	7.09	-1.74	. 00	-7.42
	KIN	. 88	-		2.04			- 3. 0 3		11
7 : 17	SNA	-1.05	-	99	-1.35			1.32	.00	-2.07
LIN	OEI	. 69	-		. 28	3.46	- 3.57	03		.83
	тот і	. 53	-		. 97	3.46	- 3, 57	-1.74	. 00	-1.34

DENSITY PARTITIONING FOR VALENCE AO'S

	Valence Atomic Orbitals		Overlap	·	ons					
	k	s	σ	S _(Li.H) _	p	v	q	q ^N	q ^T	P(Li, H)
i	. 999617	011785	025048	. 07 47 30	2.010	008	2.002	2.008	00ó	106
b	. 024203	. 811232	. 584224	. 704494	. 309	. 378	. 688	. 991	303	. 537
l	. 01 3435	584606	.811206	. 178408	. 000	. 000	.000	. 000	. 000	. 000
h		(unhybridize	d)	-	. 939	. 371	1.309	1.000	. 309	-

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

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	ĺ		Li	<u> </u>	н
		i	ъ	l	h
i	PF SPN SPT TOT	2.007 .000 .000 2.007	. 002 003 . 001 001	. 000 . 000 . 000 . 000	. 00 3 007 004
Ъ	PR SPN SPT TOT	. 002 003 . 001 001	. 989 -1. 188 . 435 . 237	.000 .000 .000 .000	1. 191 739 . 452
£	PR SPN SPT TOT	.000 .000 .000 .000	.000 .000 .000 .000	.000 .000 .000 .000	- . 000 . 000 . 000
h	PR SPN SPT TOT	.003 007 004	1.191 739 .452	. 000 . 000 . 000	1.000 -1.194 1.056 .862

Figure 8. Binding energy decomposition and description of electron distribution for LiH, SAO approximation

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	I		romotion			Quasi-	Shi Penel	ring tration	SP+QC
Orbital	×		¢	Total	Orbital	QCT	SPN	SPT	Total
L1 & NA DEI TOT	.77 -1.71 .66 23	.08 09 -3.00 -3.00	.00 .00 2.99 2.99	-16 -1.7e -72 21	KIN L1 1 OEI TOT	64 .77 .81 .14	.27 .07	21	64 .77 .26 .19
L1 = OEI TOT	.25 29 -3.00 -3.00	-2.00 9.85 .06 7.91	09 22 - 28 - 80	-1.83 9.66 -5.98 1.85	KIN SNA LI 6 OEI TOT	-1.53 .57 .36 59	4.59 4.59	-1.67 -1.67	-1.53 .57 J.23 2.32
LI J DEI TOT	2.99 2.99 2.99	99. 99 99 99	1.95 -3.01 21 -7.06	1.95 -9.21 5.95 -1.11	LI & OEI Tot	.00 .00 .00 .00	02 00	. 60 . 72	. 20 . 20 20 30
Total Li KIN NA OEI TOT				• RB -1.05 • 69 • 53	Total Li Kin SNA OEI TOT	-2.17 1.34 .37 46	4.65	-1.65 -1.65	-2.17 1.34 3.34 2.52
					Tetal H KIN SNA OEI TOT	4.21 -3.16 .82 1.87	10.15 10.15	-9.98 -3.98	4.21 -3.16 1.99 3.84

INTRA-ATOMIC ENERGY CONTRIBUTIONS OF LIN, SAU

Orbital	Orbital Quasiclassical Electrostatic Pair								Totals								
LiH		QCN		QCT		(Tellog	SPN	SPT	SP(N+T) SIN			SIT	SI(N+T)	5	T	N+T
ih (KIN SNA DEI TUT	C4 24	. ⁵⁵ .	.02	.02	86 82 88	03 03	. 27 . 27	.24 .24	.97 11 84 .73	-,18	.ee	.02 .02	.97 19 24 .75	.97 24 07 .65	22 .05 .25	.97 25 22 .71
bh (KIN SNA DEI FOT	95 95	.32	.89 .40	.17	46 89 -1.35	-11.31 -11.31	7.02 7.02	- 4. 29 - 4. 29	-4.01 .52 .01 -2.47	1.00	07	.06 .81	-4.01 1.51 .01 -2.48	-4.81 .58 -11.30 -14.73	.48 6.13 6.61	-4.21 1.05 -5.17 -6.13
2 h (KIN SNA DEI TOT	.ee .ee	.26	.00 .00	.00	. 30 . 89 . 90	.88 .22	. 80 . 20	. 22 . 22	99. 99. 99.	.00	. 00	.92 .82	. 28 . 28 . 38	.00 .20 .22 .99	.00 .00 .99	55. 59. 55. 35.
Bond H Total S	CIN SNA DEI TOT	39 99	.32	.91 .44	. 15	52 91 -1.43	-11,34	7.09 7.29	-4.25 -4.25	-3.83 -42 83 -1.74	,90	07	.27	-3.03 1.32 03 -1.74	-3.23 .33 -11.37 -14.07	.47 6.18 6.65	-3.03 .80 -5.19 -7.42

INTERATOMIC ENERGY CONTRIBUTIONS OF LIH, SAO

Figure 8 (Continued)

		Prom	otion	Qua clas	asi- sical	Sha Penet	ring ration	Sharing Interference		
		PRH_	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
	KIN	1.02	-		-2.10					-1.08
• •	NA	-1.22	-		1.33					. 11
1_1	OEI	. 70	-		. 37	4.56	-1.64			3.99
	тот	. 50	-		- 、40	4.56	-1.64			3.02
	KIN	_	-		4, 27					4.27
	SNA 1	-	-		- 3.20					- 3. 20
н	OEI	-	-		.84	10.16	-9.03			1.96
	тот	-	-		1.91	10.16	-9.03			3.03
	KIN							- 2, 95		-2.95
	SNA			96	. 43			1.25	. 01	.74
BOVD	OEI				92	-11.22	7.03	03		-5.15
	тот			96	49	-11.22	7.03	-1.72	. 01	-7.36
	KIN	1.02	-		2, 17			-2.95		. 24
	SNA	-1 22	-	96	-1.44			1.25	. 01	-2.35
LIH	OEI	. 70	-		. 29	3.50	- 3. 64	03		. 80
	TOT	. 50	-	- 96	1.02	3.50	- 3. 64	-1.72	. 01	-1.31

DENSITY PARTITIONING FOR VALENCE AO'S

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	Valence Atomic Orbitals		ne Orbitals Overlap Bon					Sondorders and Populations					
	k	8	٣	S(Li, H)	р	v	q	۹ ^N	Tp	^p (L:, H)			
i	. 999593	013973	024800	. 074058	2.010	008	2.002	2.009	007	107			
ь	. 025640	. 812773	. 582007	. 700009	. 308	. 376	. 084	. 00l	307	. 538			
R	.012075	582413	. 812803	. 179878	. 00 0	. 000	. 000	. 000	. 000	. 000			
h		(unhybridize	d)	-	. 940	. 368	1.314	1.000	. 314	-			

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

	1		Li		Н
	1	i	ь	R	h
	PR SPN	2.007	. 002 003	. 000	. 003
1	SPT TOT	.000 2.007	.001 001	.000 .000	008 004
b	PR SPN SPT TOT	.002 003 .001 001	. 989 -1. 188 . 433 . 234	.000 .000 .000 .000	1. 191 741 . 451
R	PR SPN SPT TOT	. 000 . 000 . 000 . 000	.000 .000 .000 .000	.000 .000 .000 .000	. 000 . 000 . 000
h	PR SPN SPT TOT	.003 008 004	- 1.191 741 .451	.000 .000 ·	1.000 -1.195 1.062 .868

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Figure 9. Binding energy decomposition and description of electron distribution for LiH, BAO approximation

1		Pr	rozotion —			Quasi- classical	She	ring ration	SP+Q
Orbital	×		۳	Total	Orbital	2CT	528	575	TOLA.
L1 K OEI TOT	.77 -1.72 .67 28	.14 15 -2.91 -2,93	.02 .02 2.91 2.91	.76 -1.71 .74 21	KIN L1 1 OEI TOT	64 .78 .01 .14	.06 .96	ei ei	64 .78 .06
L1 S OEI TOT	-14 -2.91 -2.93	-1.88 9.57 .PC 7.74	- 20 - 20 - 20 - 20	-1.62 9.25 -5.82 1.83	LI b OEI TOT	-1.46 .55 .36 55	4.52 4.52	-1.63 -1.63	-1.46 .55 3.23 2.52
LI J O OEI TOT	.00 .00 2.91 2.91	.00 .00 .00 .00	1.86 -5.76 01 -6.91	1.86 -5.76 5.79 -1.11	L1 & OE1 TOT	.82 .20 .02 .20	02 00	.22 .20	29. 20. 95 99
Tatal Li KIN NA OEI TOT				1.02 -1.22 .70 .58	Total Li Kin SNA OEI TOT	-2.10 1.33 .37 40	4.56 4.56	-1.64 -1.64	-2.10 1.33 3.29 2.52
					Total H KIN SNA OEI TOT	4.27 -3.28 .94 1.91	12.16	-9.23 -9.23	4.27 -3.28 1.96 3.03

INTRA-ATOMIC ENERGY CONTRIBUTIONS OF LIH, BAO

INTERA IUMIC	ENERGY	CUNTEIBUTIONS	OP	L1H.	BAO

Orbit	n)	Quasi	classical Electro	(Sharing] Totals		
Pai	r				Penetration			- 1 Interference						
LIH	-	SCH	QCT	QC(N+T)	SPN	SPT	SP(N+T)	SIN		SIT	SI(N+T)	N	T	N+T
i h	KIN SNA OEI TOT	24 24	.ep02 02 04	26 22 08	83 83	. e7 . e7	. 84 . 24	1e .97 24 .73	1e	.8e .82 .82	.97 18 24 .75	.97 24 87 .66	82	.97 24 22 .71
b h	KIN SNA GEI TOT	92 92	.31 .14 91 45	47 91 -1.37	-11,19 -11,19	6.96 6.96	-4.23 -4.23	-3.92 .46 .01 -2.45	.99	07 .05 81	-3.92 1.44 .21 -2.46	-3.92 .53 -11.13 -14.56	.44 6.25 6.52	-3.92 .97 -5.12 -8.27
l h	KIN SNA OEI TOT	.00 .00	.99 .99 .99	. 22 . An . 20	.82 .89	. 82 . 82	.00 .28	.20 .80 .20	.22	.22 52. 52.	95. 59. 55. 80.	.00 .23 .02 .20	.22 .22 .88	29. 29. 59.
Bond Total	KIN Sha Dei Tot	96	.31 .12 92 49	53 92 -1.45	-11.22	7.03 7.03	-4.19 -4.19	-2.95 .36 83 -1.72	.19	27 .08 .¢1	-2.95 1.26 23 -1.71	-2.95 .30 -11.25 -13.90	.44 6.12 6.54	-2.95 .74 -5.15 -7.36

Figure 9 (Continued)

	.	Pro	notion	Qua	asi- Isical	Sha Penet	ring ration	Sharing Interference		
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
Li	KIN NA OEI TOT	.96 -1.34 .80 .43	3.28 -5.87 3.06 -48		-3.35 1.90 .55 90	5.28 5.28	-1.76 -1.76			. 89 - 5. 30 7. 93 3. 51
H•	KIN SNA OEI TOT	. 00 . 00 . 00 . 00	63 .64 .00 .01		4.60 -3.53 1.04 2.11	9.87 9.87	-9,22 -9,22			3.97 -2.89 1.69 2.77
BOND	KIN SNA OEI TOT			-1,17 -1 17	.78 -1.25 48	-11.88 -11.88	7.56 7.56	- 3. 21 1. 51 03 - 1. 75	.00 .00	-3.21 1.13 -5.61 -7.69
LiH	KIN SNA OEI TOT	.90 -1.34 .80 .43	2.05 -5.23 3.06 49	-1,17 -1,17	1.25 85 .34 .73	3. 27 3. 27	- 3, 42 - 3, 42	-3.21 1.51 93 -1.75	.00	1.65 -7.06 4.01 -1.41

DENSITY PARTITIONING FOR VALENCE AO'S

	Valence Atomic Orbitals			Overlap	Bondorders and Populations								
	k	5	σ	S _(Li,H)	þ	v	q	qN	т _р	^p (Li, H)			
÷	999 010	009977	026031	.077872	2.011	009	2.002	2.011	008	114			
ь	.023080	.821135	. 570267	. 720442	. 271	. 372	. 043	. 989	346	. 517			
l	. 015734	570640	. 821045	. 182277	. 000	. 000	. 000	. 000	. 000	. 000			
h		(unhybridize	d)	-	. 991	. 363	1.354	1.000	. 354	-			

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

	1		Li		н
		i	ь	2	h
	PR	2.007	. 00 3	. 000	-
	SPN	. 000	005	. 000	. 005
1	SPT	. 001	. 001	. 000	009
	тот	2.008	001	.000	005
	PR	. 00 3	. 986	. 000	-
	SPN	005	-1.179	. 000	1.184
Ь	SPT	. 001	. 400	. 000	747
	TOT	001	. 207	.000	. 437
	PR	. 000	. 000	. 000	-
^	SPN	. 000	. 000	. 000	. 000
x	SPT	. 000	. 000	. 000	. 000
	тот	. 000	. 000	.000	. 000
	PR	-	-	-	1.000
	SPN	. 005	1.184	.000	-1.189
п	SPT	009	747	. 000	1.111
	TOT	005	. 437	. 000	. 922

Figure 10. Binding energy decomposition and description of electron distribution for LiH, BMAO approximation

1	Promotion								ring	SP+QC				
Orbital	k PBH	PRC	PRH	PRC	PRH	PRC	PBH	Total PBC	PB(H+C)	Orbital	classical QCT	Penet SPN	ration SPT	Total
KIN NA Li k DEI TOT	.94 -2.11 .81 36	.65 72 .06 82	.02 03 -2.81 -2.82	.00 01 .62 .62	.88 .89 2.79 2.79	.80 .80 .89 .89	.94 -2.11 .89 27	.65 72 .89 .01	1.59 -2.83 .98 26	KIN SNA Li i OEI TOT	81 .97 .02 .17	.09 .09	02 02	81 .97 .49 .25
KIN NA DEI TOT	.02 03 -2.81 -2.82	.00 01 .62 .62	-1.82 9.23 .86 7.47	1.28 -2.40 01 -1.13	.00 .00 00 00	.00 .00 02 00	-1.77 9.17 -5.63 1.77	1.28 -2.41 1.22 .10	48 6.77 -4.41 1.87	L1 b SHA DEI TOT	-2.54 .93 .54 -1.08	5.19 5.19	-1.75	-2.54 .93 3.98 2.36
KIN LI G OEI TOT	.02 .20 2.79 2.79	.00 .00 .89 .89	.00 .00 00 00	.00 .00 00 00	1.78 -8.40 01 -6.62	1.35 -2.74 61 -1.39	1.78 -8.40 5.55 -1.07	1.35 -2.74 1.75 .37	3.14 -11.14 7.30 70	LI L SNA LI L OEI TOT	. 00 . 00 . 00 . 00	00 00	. 88 • 88	.00 08 80
Total Li KIN NA OBI TOT							.96 -1,34 .80 .43	3.28 -5.87 3.06 .48	4.24 -7.20 3.86 .98	Total Li KIN SNA UEI TOT	-5.35 1.90 .55 90	5,28 5,28	-1.76 -1.76	-3.35 1.90 4.67 2.61
Total H KIN HA OEI TOT			• ······•				.00 .99 .00 .99	63 .64 .00 .01	63 .64 .80 .91	Total H KIN SNA OEI TOT	4.60 -3.53 1.04 2.11	9.87 9.87	-9.22 -9.22	4.68 -3.53 1.69 2.76

INTRA-ATOMIC ENERGY CONTRIBUTIONS OF LiH, BMAO

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INTERATOMIC ENERGY CONTRIBUTIONS OF LIN, BMAO

Orbit	1	Qui	asiclassical Electro	static			Sh	aring					-Totals -	· ·
Pal: Li H	r .	QCN	QCT	QC(N+T)	SPN	Penetratic SiT	SP(N+T)	SIN	Int	erference SIT	51(N+T)	×	T	X+T
ih	KIN SNA OEI TOT	04	.0003 05	07 03 10	04 04	. <i>0</i> 8 .08	. 84 . 84	14 85 .77	•11	.01 .02 .03	1.06 ~.22 ~.05 .80	1.86 29 89 .69	08 .06 .85	1.86 29 03 .75
bh	KIN SNA OEI TOT	-1.13	.43 .38 -1.23 42	32 -1.23 -1.55	~11.84 ~11.84	7.47 7.47	-4.37 -4.37	-4.27 .89 .01 -2.58	•B7 -	.07 .05 02	-4.27 1.74 .01 -2.52	-4.27 .63 -11.83 -15.47	.79 6.25 7.83	-4.27 J.42 -5.58 -8.44
£ h	kin Sna Oei Tot	.00 .00	00, 00, 99, 99,	. 00 . 89 . 80	.00	. 08 . 00	. 88 . 88	.08 .80 .80 .89	.00	88. 88. 88.	. 89 . 98 . 98 . 99	.00 .00 .00 .00	. 80 . 80 . 80	. 88 . 88 . 88 . 88
Bond Total	KIN SNA OEI TOT	-1.17 -1.17	•43 •35 •1.25 48	-,39 -1,25 -1,64	-11.88	7.56	-4.33 -4.33	-3.21 .75 83 -1.73	.76 -	.07 .07 .67	-3.21 1.52 83 -1.72	-3.21 .35 -11.92 -14.78	.78 6.30 7.89	-3,21 1,13 -5,61 -7,69

Figure 10 (Continued)

	l	Prom	otion	Qua	asi- sical	Sha Pene	ring	Sha: Interi	ring	
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
в	KIN NA OEI TOT	62 . 40 1. 43 1. 22	- - -		-2.33 1.92 .10 31	4. 20 4. 20	. 18 . 18			- 2.95 2.32 5.91 5.29
н	KIN SNA OEI TOT	- - -	- - -		1.65 -1.24 .12 .54	4. 42 4. 42	-1.89 -1.89			1.65 - 1.24 2.65 3.06
BOND	KIN SNA OEI TOT			-1.85 -1.85	30 18 54	-6.60 -0.00	. 85 . 85	-4.00 1.97 14 -2.17	. 13	- 4.00 11 - 6.06 -10 17
вн	KIN SNA OEI TOT	ɛ.: . 40 1. 43 1. 22	-	-1.85	08 . 32 04 31	2.02	80 80	-4.00 1.97 14 -2.17	. 13	- 5.30 .97 2.50 - 1.82

DENSITY PARTITIONING FOR VALENCE AO'S

	Valence Atomic Orbitals			Cvurlay	Bondorders and Populations								
	k	*	œ	⁵ (В, Н)	р	x	q	q ^N	q ^T	^P (B, H)			
:	. 277088	178357	. 009140	. 000000	2.000	. 000	2.000	2 000	. 000	. 000			
2	. 209437	. 895130	393540	. 313600	2.190	141	2.049	2.080	032	451			
5	.010173	. 399245	. 916701	. 702931	. 428	. 403	. 830	. 920	090	. 573			
h		(unhybridize	d)		. 860	. 201	1.121	1.000	. 121	-			

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

1		— В —		н
	i	2	b	h
FR	2.000	. 000	. 000	-
SPN	. 000	. 000	.000	. 000
SPT	, 000	. 000	.000	. 000
TOT	2.000	. 000	. 000	. 000
PR	000	2 080	.000	-
SPN	000	- 001	- 018	. 019
SPT			004	- 096
		2 140	- 014	- 077
101		2.140	014	017
PR	. 000	. 000	. 920	-
SPN	. 000	018	483	. 501
SPT	. 000	.004	087	006
TOT	. 000	014	. 350	. 495
PR	-	-		1.000
SPN	. 000	. 019	. 501	520
SPT	. 000	096	006	. 223
тот	. 000	077	. 495	. 703
	FR SPN SPT TOT PR SPN SPT TOT PR SPN SPT TOT PR SPN TOT	i FR 2.000 SPN .000 SPT .000 TOT 2.000 PR .000 SPN .000 SPT .000 SPT .000 SPT .000 SPT .000 SPN .000 SPT .000 SPT .000 SPT .000 SPT .000 TOT .000 SPT .000 SPT .000 TOT .000	I B i \$\$ FR 2.000 .000 SPN .000 .000 SPT .000 .000 TOT 2.000 .000 SPN .000 .000 TOT 2.000 .000 PR .000 .001 SPT .000 .001 SPT .000 .000 TOT .000 .000 SPN .000 .014 PR .000 .014 PR .000 .019 SPT .000 .019 SPT .000 .017	i g b i g b FR 2.000 .000 .000 SPN .000 .000 .000 SPT .000 .000 .000 TOT 2.000 .000 .000 PR .000 2.080 .000 SPN .000 .001 .018 SPT .000 .000 .004 TOT .000 .000 .920 SPN .000 .018 .483 SPT .000 .004 .087 TOT .000 .014 .350 PR .000 .014 .350 PR .000 .019 .501 SPT .000 .019 .501 SPT .000 .096 .006 TOT .000 .096 .006

Binding energy decomposition and description of electron distribution for BH, SAO approximation Figure 11.

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1		P			1	Quasi-	Sha	ring	SP+QC
Orbital	k	•	~	Total	Orbital	QCT	SPN	SPT	Total
KIN B k OEI TOT	.97 -2.07 .01 -1.09	62 .63 -1.71 -1.67	.02 .22 1.91 1.91	1, 01 -2,11 ,17 -,93	B 1 OEI TOT	93. 69. 99.	- 36	.00	53. 22. 53 93
KIN B S OEI TOT	60 -63 -1.71 -1.67	-2.81 10.46 03 7.62	.20 .20 96 96	-4.84 11.77 -1.14 6.59	KIN B & SNA B & OEI TOT	31 .44 .02 .16	. 35 . 35	-,46 -,46	31 .44 63 .65
B J CEI TOT	.00 .00 1.91 1.91	.20 .00 96 96	2.41 -9.26 2.97 -3.88	2.41 -9.26 2.40 -4.45	B b SNA B b OEI TOT	-2.03 1.48 .CB 47	3.85 3.85	.63 .63	-2.03 1.45 4.56 4.22
Total B XIN NA OEI TOT				62 .40 1.43 1.22	Total B KIN SNA OEI TOT	-2.33 1.92 .10 31	4.20 4.20	,15 ,18	-2.33 1.92 4.49 4.07
					Total H KIN SNA OEI TOT	1.65 -1.24 .12 .54	4.42	-1,89	1.65 -1.24 2.65 3.06

INTEA-ATOMIC ENERGY CONTRIBUTIONS OF BH, SAG

INTERATURIC	ENERGY	CONTRIBUTIONS	ÚΡ	BH.	SAO

Orbital) Qua	siclassical Electro	ostatic	·		Sh	aring					- Totals	
B N	QCN	2CT	QC(N+T)	SPN	SPT	SP(N+T)	SIN	Interre	SIT	SI(S+T)	N	T	N+T
i h OEI TOT	84	.8983 .88 83	07 .02 07	- 88 - 89	. C2 . 00	.02 .00	.30 12 18	0e .eo	.00	.00 .30 12 .18	.82 .25 12 .14	03 .00 23	.02 .23 12 .11
L h SHA DEI TOT	.45 .45	.0282 23 84	36 03 39	84 84	:77	.73 .73	1.2797 07 3,36	98 .82	.e3 .85	3.26 .21 27 3.48	3.26 .62 11 3.77	76 .74 02	3.26 14 .63 3.75
KIN Sha Dh Qei Tot	-2.26	.25 .22 14 .33	-1.79 14 -1.93	-6.56 -6.56	. 28 . 28	-6.47 -6.47	-7,26 2,01: -5,70	5eei	.10 .09	-7.26 1.60 .25 -5.62	-7.26 75 -6.51 -14.52	.56 86 .58	-7.26 -19 -6.57 -14.83
Bond KIN Total SNA OEI TOT	-1.85	.2662 18 54	-2.22 18 -2.39	-6.68	.85	-5.74	-4.00 3.37 -1.4 -2.17	40 .00	.13	-4.80 2.11 14 -2.03	-4.00 .12 -6.74 -10.62	23 .68 .45	-4.02 11 -6.05 -12.17

Figure 11 (Continued)

		Prom	notion	Qua	asi- sical	Sha Pene	ring tration	Shar Interf	ing erence	
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
в	KIN NA OEI TOT	-1.13 1.71 .65 1.23	- - - -	<u></u>	-1, 97 1, 06 .08 23	4. 18 4. 18	. 04 . 04	- <u></u>		- 3.10 3.37 4.95 5.22
н	KIN SNA OEI TOT		- - -		1.47 -1.10 .10 47	4.59 4.59	-1.80 -1.80			1.47 - 1.10 2 89 3 26
BOND	KIN SNA OEI TOT			-1 ->0 -1, >0	- 31 14 45	-7 02 -7 02	1 10 1.1 0	-4.07 2.26 - 13 -1.74	. 11	- 4.07 .15 - 6.19 -10.10
вн	KIN SNA OEI TOT	-1.13 1.71 .65 1.23	- - -	-1. 90 -1. 90	50 .25 .04 21	1.75 1.75	00 00	-4.07 2.20 13 -1 34	. 11	- 5.70 2.42 1.65 -1.62

DENSITY PARTITIONING FOP VALENCE AO'S

	Valence Atomic Orbitals			Overlap	Bondorders and Populations						
	k	\$	÷	⁵ (В.Н)	p	v	q	q ^N	ч ^Т	^р (В.Н)	
1	. 972426	217807	. 083358	. 000000	2.000	. 000	2.000	2.000	. 000	. 000	
L	. 232611	. 880182	413729	. 281115	2.148	109	2.039	2.062	023	380	
ь	.016743	. 421711	.900576	.727660	. 440	. 413	. 853	. 238	086	. 568	
h	1	(unhybridized	d)	-	. 804	. 305	1.108	1.000	. 108	-	

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

	i	ВВ			н
		i	£	ь	h
	PR SPN	2.000	. 000	. 000	. 000
1	SPT TOT	. 000 2. 000	. 000 . 000	. 000 . 000	. 000 . 000
	PR	. 000	2.085	023	-
l	SPT TOT	. 000 . 000 . 000	. 026 2. 110	.003	051 051
b	PR SPN SPT TOT	. 000 . 000 . 000 . 000	023 .009 .003 011	. 961 558 036 . 368	- .549 052 .496
h	PR SPN SPT TOT	- . 000 . 000 . 000	009 051 060	- .549 052 .496	1.000 540 .211 .672

Figure 12. Binding energy decomposition and description of electron distribution for BH, BAO approximation

}			Promotion		1	Quasi- classical	Sharing Penetration		SP+QC	
Orbital	k _	t	σ .	Total	Orbital	QCT	SPN	SPT	Total	
KIN B k OSI TOT	.98 -1.93 .02 -1.81	40 .42 -2.36 -2.33	.00 .00 2.33 2.33	.92 -1.95 .17 86	B 1 OEI TOT	.00 .20 .30 .P0	00 00	. ee . 55 .	68. 96. 69 52	
HIN B = NA DEI TOT	48 .42 -2.36 -2.33	-3,99 14.16 55 9.62	.00 92 92	-4.80 15.03 -2.96 7.27	B P OEI TOT	22 .35 .01 .15	.13 .13	20 22	22 .35 06 .07	
B or OEI TOT	.00 2,33 2,33	.20 .00 92 92	2.75 -11.36 3.09 -5.53	2.75 -11.36 3.43 -5.18	B b OEI TOT	-1.75 1.31 .26 38	4.26 4.26	.24 .24	-1,75 1,31 4,36 3,92	
Total B KIN NA OBI TOT				-1.13 1.71 .65 1.23	Total B KIN SNA OEI TOT	-1.97 1.66 .08 23	4.18 4.15	.04 .24	-1.97 1.66 4.32 3.99	
					Total H KIN SNA OEI TOT	1.47 -1.10 .12 .47	4,59	-1,50 -1,80	1.47 -1.18 2.89 3.26	

INTRA-ATOMIC ENERGY CONTRIBUTIONS OF BH, BAO

INTERATORIC	ENERGY	CONTRIBUTIONS	G.P	BH .	940
ANICANIUNIC	CUCUOI	CONTRIBUTIONS	OF.	on,	DAU

Orbita	1	Quasi	Quasiclassical Electrostatic									Totals		
B R		QCN	QCT	QC(N+T)	SPN	SPT	SP(N+T)	SIN	SIT	SI(N+T)	8	T	N+T	
i h	KIN SNA OEI TOT	24 24	.ec03 .ee e3	87 .20 87	. 90 . 80	.02 .80	.88 .82	.80 .31 .02 18 .22	.20 .90 99.	.22 .31 12 .22	.20 .27 10 .17	23 .22 23	.00 .24 10 .14	
L h	KIN SNA OEI TOT	.45	.2973 02 74	24 62 26	.16 .16	. 41 . 41	.57 .57	2.66 1.3079 07 3.10	.02 .03 .95	2.66 .56 87 3.15	2,66 .99 .09 3.74	• 67 • 39 • 28	2.66 .32 .48 3.46	
bh	KIN SNA DEI TOT	-2.34	.25 .20 12 .32	-1.90 12 -2.02	-7.18 -7.18	.69 .69	-6.49 -6.49	-6.73 1.8236 -5.25	82 .88 .Ø€	-6.73 1.49 .24 -5.15	-6.73 91 -7.14 -14.77	.50 .57 1.07	-6.73 41 -6.57 -13.72	
Bond Total	KIN SNA DEI TOT	-1.98	.2556 14 45	-2.21 14 -2.35	-7.82 -7.82	1.1e 1.1e	-5,92 -5,92	-4.07 3.41 -1.15 13 -1.94	23 .11 .11	-4.27 2.36 13 -1.83	-4.27 .36 -7.15 -12.86	28 .96 .76	-4,07 ,15 -5,19 -18,10	

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Figure 12 (Continued)

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		Promotion		luasi- classical		Sharing Penetration		Sharing Interierence		
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
В	KIN NA OEI TOT	30 1 97 . 25 1. 42	06 -1.57 1.77 .12		50 .33 .30 - 10	3, 27 3, 77	16 16	· .		- 1, 36 .77 6, 15 5, 57
H	KIN SNA OEI TOT	. 00 . 00 . 00 . 00	5, 53 -5 06 , 00 47		. 45 2 + . 01 . 17	4. 09 4. 09	- 27 - 27			5,98 - 5,35 4,43 5,06
EOND	KIN SNA OEI TOT			-1.70 -1.70	07 91 96	- 6, 60 - 6, 60	. 00 . 00	-6.91 2.62 08 -4.37	. 02 . 02	- 6.91 .80 - 6.69 -12.79
BH	KIN SNA OEI TOT	- . 80 1 97 25 1. 42	5.47 -0.05 1.77 53	-1.70 -1.70	05 .03 00 02	2.06	- 11 - 11	-6.91 2.62 08 -4.37	. 02	- 2,29 - 3 78 3,89 -2,16

DENSITY PARTITIONING FOR VALENCE AO'S

	call all Atomic Orbitals			Overlap	rlap Bondorders and Populations						
	÷.		σ	^S (В. Н)	þ	v	q	q ^N	qT	^р (В. Н)	
:	. 27 37 27	212751	. 081189	00000	2.000	. 000	2.000	2.000	. 00 0	. 00 0	
L	227053	. 880058	417105	.239642	2.103	074	2.029	2,032	004	309	
ь	. 017290	. 424581	. 205225	. 686640	. 542	. 406	. 948	. 908	020	. 592	
h		(unhybridize	d)	-	692	. 332	1 024	1.000	. 024	-	

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

	ļ		В					
		i	e	ь	h			
i	PP SPN SPT TOT	2.000 .000 .000 2.000	. 000 . 000 . 000 . 000	.000 .000 .000 .000	- . 000 . 000 . 000			
1	PR SPN SPT TOT	. 000 . 000 . 000 . 000	2.077 002 .003 2.077	044 .036 .000 008	034 007 041			
þ	PR SPN SPT TOT	.000 .000 .000 .000	044 .036 .000 008	1.012 536 024 .452	- . 499 . 004 . 503			
h	PR SPN SPT TOT	- 000 000 000	034 007 041	- 499 - 004 - 503	1.000 465 .027 .561			

Figure 13. Binding energy decomposition and description of electron distribution for BH, BMAO approximation

	1				Promoti	on		Total		1	Quasi-	Sha	ring	SP+QC
Orbital	ран	K PBC	PBH	PEC	PBH	PBC	PBH	PHC	PR(H+C)	Orbital	QCT	SPN	SPT	Total
B & OEI TOT	.40 86 02 48	.28 30 .04 .02	.09 09 -2.76 -2.77	90 .00 -1.02 -1.82	.80 .98 2.66 2.66	.00 .00 1.62 1.62	.40 86 .85 41	- 39 - 86 - 84	-1.16 -11 37	KIN B 1 OEI TOT	. 82 . 88 . 80 . 86	-,00 -,00	.00	69 98 98
KIN NA B = ORI TOT	.09 09 -2.76 -2.77	00 .00 -1.02 -1.02	-4.56 16.17 -1.16 10.46	-4.57 7.09 32 2.20	.00 68 68	.00 .00 .31 .31	-4.38 15.99 -4.96 6.65	-4.57 7.09 -2.99 38	-8.96 23.07 -7.85 6.27	B & SNA B & OEI TOT	84 .06 .8P .02	00 00	02 02	84 .46 83 81
KIN B of OEI Tot	.00 .00 2.66 2.66	.00 .00 1.62 1.62	.00 .00 ~.68 ~.68	.00 .00 .31 .31	3.18 -13.15 2.97 -7,00	4.24 -8.38 .22 -3.93	3.18 -13.15 5.15 -4.52	4.24 -8.39 4.60 .45	7.42 -21.53 9.75 -4.36	KIN SNA B b OEI TOT	46 .34 .88 12	3.97 3.97	. 18 . 18	46 .34 4.16 4.24
Total B KIN H/ QEI TOT							+.90 1.97 ,25 1.42	06 -1.59 1.77 .12	96 .38 2.01 1.74	Total B KIN SNA OEI TOT	50 .39 .00 10	3.97 3.97	.16	58 .39 4.14 4.85
Total H KII NI DE TO							. 98 . 99 . 98 . 98 . 93	5.53 -5.06 .00 .47	5.53 -5.86 .80 .47	Total H KIN SNA OEI TOT	.45 29 .01 .17	4.69	27 27	.45 29 4.43 4.59

INTRA-ATONIC ENERGY CONTRIBUTIONS OF BH, BMAO

INTERATOMIC ENERGY CONTRIBUTIONS OF BH, BMAD

Orbital	Quasi	Quasiclassical Electrostatic						ring				Totale			
Pair B H		QCN	QCT	QC(X+T)	SPN	SPT	SP(K+T)	SIN		interfe	SIT	SI(N+T)	N	T	N+T
1 h	KIN SNA OEI TOT	03	.0801 .09 01	84 .08 84	.00 .00	. es . vo	.00 .00	.19 07 .12	ិពធ	, 643 , 1	.#E	.88 .19 87 .12	.80 .16 07 .89	01 .00 01	.88 .15 87 .89
Lh	KIN SNA OBI TOT	.22 .22	.8817 80 17	.05 09 .05	. 35 . 35	. 86 . 86	.41 .41	2,42 .60 05 2,34	63	. <i>RO</i>	. <i>01</i>	2.42 02 05 2.34	2.42 .19 .29 2.90	16 .06 10	2.42 .83 .35 2.90
bh	KIN SNA OEI TOT	-1.94	.04 .06 01 .10	-1.84 01 -1.85	-6.95 -6.95	06 06	-7.01 -7.01	-9.33 2.12 .6.83	.33	•.00	.01 91	-9.33 2.46 .04 -6.32	-9.33 .58 -6.91 -15.73	.11 06 .05	-9.33 .62 -6.97 -15.68
Bond Total	KIN SNA OEI TOT	-1.76	.0411 01 08	-1.83 01 -1.84	-6.68 -6.60	.00 .00	-6.68 -6.68	-6.91 2.92 08 -4.37	-,30	00	.02	-6.91 2.63 08 -4.36	-6.91 .85 -6.68 -12.73	85 81 86	-6.91 .90 -6.69 -12.79

Figure 13 (Continued)

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Figure 14. Binding energy decomposition and description of electron distribution for NH', SAO approximation

		Promotion		Quasi- classical		Sharing Penetration		Sharing Interierence			
	1	PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL	
N	KIN NA OEI TOT	- 5.13 17.24 7 97 20 07			27.22 -26.86 2.06 2.42	12 26 12 26	-12.71 -12.71	<u>. </u>		22,07 - 9,62 9 58 22,04	
н	KIN SNA OEI TOT	۲ <u>-</u> - -			- 6.26 4 70 1.80 24	10.62 10.62	- 3.30 - 3.30			- 6.26 4.70 9.06 7.49	
BOND	KIN SNA OEI TOT			-1 17 -1.17	1.38 - 2.80 - 1.43	- 16. 51 - 16. 51	11. 31 11. 31	-14.93 7.12 31 - 8.12	. 12	-14.93 7.44 - 8.30 -15.80	
NН	KIN SNA OEI TOT	- 5,13 17 24 7,97 20,07	- - -	-1.17 -1.17	20, 96 -20, 78 1, 06 1, 23	6.37 6.37	- 4.76 - 4.76	-14 93 7.12 - 31 - 8 12	. 12	.90 2.52 10.34 13.73	

DENSITY PARTITIONING FOR VALENCE AO'S

.

	al ace Atomic Orbitals			Overlap	rlap Bondorders and Populations						
	ĸ	5	س	^S (N. H)	þ	v	ą	q.N	$\mathbf{T}_{\mathbf{p}}$	^р (N, H)	
i	. 999821	0182c7	004915	. 052271	2.004	002	2,001	1. 228	. 004	045	
ь	.018863	. 982294	. 186394	. 590041	1.153	. 307	1.459	1.002	. 457	. 520	
2	.001423	180453	. 982465	. 292040	. 000	. 000	. 000	. 000	. 000	. 000	
7		(unhybridize	d)	0	2.000	. 000	2.000	2.000	.000	. 000	
h		(unhybridize	d)	-	. 235	. 304	. 540	1.000	460	-	

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

i.

		·	N					
		i	b	g	fr	h		
	PR	2.003	005	. 000	. 000			
	SPN	. 000	. 005	. 000	. 000	005		
1	SPT	. 000	. 000	. 000	. 000	. 004		
	тот	2.003	. 000	.000	. 000	001		
	PR	005	1.008	. 000	. 000	-		
	SPN	. 005	-1.260	.000	. 000	1.255		
b	SPT	. 000	1.317	. 000	. 000	860		
	TOT	. 000	1.065	. 000	. 000	. 395		
	PR	. 000	. 000	. 000	. 000	-		
	SPN	. 000	. 000	. 000	. 000	.000		
X	SPT	. 000	. 000	. 000	. 000	. 000		
	тот	. 000	. 000	.000	. 000	. 000		
	PR	. 000	. 000	. 000	2.000	-		
	SPN	. 000	. 000	. 000	. 000	. 000		
n	SPT	. 000	. 000	. 000	. 000	. 000		
	тот	. 000	. 000	. 000	2.000	. 000		
	PR	-	-	-	-	1.000		
	SPN	005	1.255	. 000	. 000	-1.249		
h	SPT	. 004	860	. 000	. 000	. 396		
	TOT	001	. 395	. 000	. 000	. 146		

				Promotion			1	Quasi- classical	- Sharing al Penetration SPM SPT		SP+QC Total	
Orbital	k k	8	a	r.	Ť	Total	Orbital	QCT	SPN	SPT	Total	
KIN N K OEI TOT	-1.56 3.25 31 1.38	2.92 -3.05 -26.79 -26.92	.00 .00 -25.29 -25.29	.80 .80 26.14 26.14	.00 .20 26.14 26.14	-1.54 3.24 65 1.65	KIN SNA N 1 OEI TOT	2.18 -3.26 .02 -1.06	.11 .11	.01	2.18 -3.26 .14 94	
KIN HA N s Obi Tot	2.92 -3.65 -26.79 -26.92	-62.94 212.23 -17.64 131.66	.08 .20 -17.96 -17.96	.00 .00 72 72	.00 .08 72 72	-57.12 206.15 -150.05 -1.01	KIN SNA N b OEI TOT	25.84 -23.60 2.84 3.48	12.16	-12.72 -12.72	25.04 -23.60 1.48 2.91	
KIN N C OEI TOT	.00 -25.29 -25.29	.80 .00 -17.96 -17.96	-49.92 179.20 1.70 138.98	.00 .00 -8.35 -6.35	.00 .00 -8.35 -6.35	-49.92 179.20 -126.98 2.36	N Q OEI TOT	.00 .00 .00	00 00	. 20 . 00	90. 98 99 - 99 - 92 -	
N TOFI TOT	.80 26.14 26.14	.20 .60 72 72	.00 .00 -8.35 -8.35	51.72 -185.68 28.49 -113.46	.00 .00 27.51 27.51	51.72 -185.68 142.82 8.87	N T OEI TOT	00. 99. 90.	00 00	.00	90. 90. 90 90	
NA N TOT	.88 .00 26.14 26.14	.00 .98 72 72	.00 .00 -8.35 -8.35	.00 .00 27.51 27.51	51.72 -185.68 20.49 -113.46	51.72 -185.68 142.82 8.87	N 7 SHA OBI TOT	.00 .00 .00 .00	00 00	. ce . aa	88 88 - 88 - 88 - 88	
Total N KIN NA OBI TOT						-5,13 17,24 7,97 20,07	Total H KIN SNA OEI TOT	27.22 -26.86 2.06 2.42	12.26	-12.71 -12.71	27.22 -26.86 1.61 1.97	
							Total H KIN SNA OEI TOT	-6.26 4.78 1.88 .24	10.62	-3.36 -3.36	-6.26 4.78 9.26 7.49	

INTERAATOMIC ENERGY CONTRIBUTIONS OF NH*, SAD

INTERATORIC ENERGY CONTRIBUTIONS OF MAY, SA	INTERATORIC	ENERGY	CONTRIBUTIONS	OP	ЖН,	SAO
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Orbit	Orbital Quasiclassical Electrostatic												Totals		
Pari H H	r	QCM	QCT	QC(H+T)	SPW	SPT	SP(E+T)	SIN		SIT	SI(N+T)	N	T	N+T	
i b	KIN SNA OBI TOT	84 84	-,88 .82 82 -,88	82 82 85	. 87 . 97	85 -,85	.02 .02	1.58 -,82 81 1.44	-,03	02 -,00 ø2	1.50 06 01 1.42	1.58 89 .26 1.47	.00 07 07	1.50 29 01 1.40	
b h	kin Sna Ori Tot	-1.32	9611 +2.78 +3.85	-2.39 -2.78 -5.17	-16.58 -16.58	11.36 11.36	-5.21 -5.21	-16.43 10.10 .03 -6.02	.27	.2128 .13	-16,43 12,51 .83 -5,89	-16.43 9.85 -16.54 -23.92	93 8.58 7.65	-16.43 8.12 -7.96 -16.27	
l h	kin Sna Obi Tot	.00 .00	.80 .80 .80 .90	.00 .00 .00	. 66 . 66	. 00 . 00	. 80 . 22	09. 00. 09.	. 80	.00. 99. 88.	. 88 . 88 . 99 . 88	.20 .00 .00 .00	. 60 . 20 . 22	88. 80. 98. 59.	
# h	KIN SNA ORI TQT	.18 ,19	.00 1.21 .50 1.21	1.30 .00 1.30	. 00 . 00	. ee . ee	. 88 . 88	.00 -1.60 -17 -1.77	.88	.88 .89. 6 0.	.80 -1.60 17 -1.77	.80 -1.51 -1.7 -1.67	1.21 .00 1.21	.00 30 17 46	
₩ P	KIN SNA OBI TOT	01. 01.	.00 i.21 .00 i.21	1.30 .20 1.30	.20 .20	. 80 . 80	.90 .90	.02 -1.60 17 -1.77	.ee	99. 93. 88.	.00 -1.60 17 -1.77	.00 -1.51 17 -1.67	1.21 .00 1.21	- 80 30 17 46	
Bond Total	KIN SNA OBI TOT	-3,17 -1,17	-,96 2,34 -2,80 -1,43	.20 -2.89 -2.60	-16.51 -16.51	11.31	-5.19 -5.19	-14.93 6.88 31 -8.12	.24	.2808	-14.93 7.24 31 -6.08	-14.93 5.95 -16.81 -25.80	1.49 8.51 10,00	-14.93 7.44 -8.32 -15.80	

Figure 14 (Continued)

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Figure 15. Binding energy decomposition and description of electron distribution for NH', BAO approximation

		Prom	otion	Qu cla	asi- ssical	Sha Penet	ring ration	Sha: Interfe	ring erence	
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
	KIN	- 3.82	-		27.98					24.16
N	NA	16.29	-		-28.02					-11 73
14	OEI	7.10	-		2.42	11.92	-12.99			8.46
	тот	17, 57	-		2.38	11.92	-12.99			20.89
	KIN	-	-		- 6 84					- 6,84
	SNA	-	-		5.13					5.13
н	OEI	-	-		2.15	10.48	- 3.04			9.59
	тот	-	-		. 44	10.48	- 3.04			7.88
	KIN						-	13.64		-13.64
	SNA			-1.18	1 56			6.30	.12	6.81
BOND	OEI				- 3.34	-16 28	11 36 -	. 27		~ 8.53
	тот			-1.18	- 1 77	-16.28	11.36 -	7.60	.12	-15,36
	KIN	- 3.82	-		21,14		-	13.64		3,68
	SNA	10.23	-	-1.18	-21 33			6.30	. 12	. 21
NH	OEL	7.10	-		1 23	6.12	- 4,67 -	. 27		9.52
	TOT	1 > 57	-	-1 18	1 05	6 12	- 4 07 -	7.60	.12	13.41

DENSITY PARTITIONING FOR VALENCE AO'S

	Valence Atomic Orbitals			Overlap	Overlap Bondorders and Population					
	k	2	F	S _(N, H)	р	v	q	q ^N	q ^T	^p (N, H)
i	. ???73?	022175	005484	. 050320	2,003	002	2.001	1.998	. 004	041
ь	.022811	. 981812	. 188480	. 597723	1.206	. 296	1.502	1.002	. 499	. 495
£	.001205	-,188556	. 982062	. 297275	.000	. 000	.000	.000	.000	. 000
n		(unhybridized	d)	0	2.000	. 000	2.000	2.000	. 000	. 000
h		(unhybridized	4)	-	. 204	. 294	. 497	1.000	503	-

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

)	N			
		i	ь	Q	Ħ	h
	PR	2.003	005	. 000	. 000	-
	SPN	. 000	. 005	. 000	. 000	005
1	SPT	. 000	. 000	. 000	. 000	. 004
	тот	2.003	. 000	. 000	. 000	001
	PR	005	1.007	. 0 0 0	. 000	-
	SPN	. 005	-1.243	. 000	. 000	1.238
ъ	SPT	. 000	1.363	. 000	. 000	864
	тот	. 000	1.127	. 000	.000	. 374
	PR	. 000	. 000	.000	. 000	-
•	SPN	. 000	. 000	. 000	. 000	. 000
X	SPT	. 000	. 000	.000	. 000	. 000
	TOT	. 000	. 000	.000	.000	.000
	PR	. 000	. 000	. 000	2.000	-
	SPN	. 000	. 000	.000	. 000	. 000
Ħ	SPT	. 000	. 000	.000	.000	. 000
	тот	. 000	. 000	.000	2.000	. 000
	PR	_	-	-	-	1.000
	SPN	005	1.238	. 000	. 000	-1.233
n	SPT	. 004	864	. 000	.000	. 357
	TOT	001	. 374	. 000	.000	. 124

INTEA-ATOMIC	ENERGY	CONTRIBUTIONS	OF	SH',	BAO

	!			Promotion			N	Quasi-	Sh. Penel	aring tration	SP+QC
Orbital	k		σ	n	fr	Total	Orbital	QCT	SPN	SPT	Total
KIN N K OEI TOT	-1.68 3.53 30 1.55	3.48 -3.65 -26.44 -26.62	.00 .00 -24.85 -24.85	.02 .02 25.70 25.70	.20 .00 25.70 25.70	~1.66 3.51 66 1.19	KIN SNA N 1 OEI TOT	2.13 -3.22 .02 -1.05	. 89 . 89	.01 .01	2.13 -3.20 .12 95
N S OEI TOT	3.48 -3.65 -26.44 -26.62	-68.86 269.68 -17.53 130.71	.00 .00 -17.70 -17.70	.80 .00 77 77	.00 20 77 77	-53.93 201.82 -152.69 -4.81	N D OEI TOT	25.84 -24.81 2.42 3.43	11.94 11.94	-13.00 -13.00	25.84 -24.91 1.24 2.27
KIN N O DEI TOT	.03 .00 -24.85 -24.55	.23 .82 -17.70 -17.70	-48.21 176.02 1.67 129.49	.20 .00 -8.19 -8.19	.00 .00 -8.19 -8.19	-48.21 176.82 -124.80 3.82	N L OSI TOT	.20 .20 .ec .08	- 20 - 80	. 00 . 02	98. 99. 99 99
KIN B T OEI TOT	.00 .00 25.70 25.70	.00 .00 ~.77 ~.77	.20 .20 -8.19 -8.19	49.99 -182.53 20.12 -112.44	.20 .22 27.24 27.04	49.99 -182.53 142.63 12.89	N + SNA OEI TOT	.20 .00 .02	20 22	.00 .00	89 69 - 69 - 89
N TT OEI TOT	.00 .00 25.70 25.70	.00 .00 77 77	.00 .02 -8.19 -5,19	.00 .00 27.04 27.04	49.99 -182.53 22.10 -112.44	49.99 -182.53 142.63 18.29	KIN N 77 SNA DEI TOT	55. 99. 50. 50.	20 00	. 80 . 80	89. 92. 93 99
Total N KIN NA OEI TOT						-3.82 16.29 7.10 19.57	Total N KIN SNA OEI TOT	27.98 -28.32 2.42 2.38	11.92	-12.99 -12.99	27.98 -28.02 1.36 1.32
							Total H KIN SNA OEI TOT	-6.94 5.13 2.15 .44	12.43 12.43	-3.24 -3.24	-6.84 5.13 9.59 7.86

INTERATORIC ENERGY CONTRIBUTIONS OF NR', BAO

Orbit	1	Qua	siclassical Electro	static	tic Sharing				- <u></u>	- Totals				
Pal N H	r	QCH	QCT	QC(E+T)	SPN	SPT SPT	SP(M+T)	SIN		SIT	SI(#+T)	¥	T	N+T
1 1	KIN SNA OEI TOT	05 05	-, <i>20</i> .83 82 . <i>20</i>	82 82 84	- 07 - 27	25	. 82 . 82	.29 .29 01 1.35	22	8282 82	1.29 .04 01 1.33	1.29 .02 .06 1.37	.00 07 07	1.29 .82 81 1.38
6 h	kin Sha Obi Tot	-1.34	-1.0611 -3.31 -4.49	-2.52 -3.31 -5.83	-16.35 -16.35	11.41 11.41	-4.94 -4.94	-14.93 9.04 .83 -5.62	.26	.2208 .14	-14.93 9.44 .83 -5.47	-14.93 7.96 -16.32 -25.29	-1.04 8.09 7.05	-14.93 6.92 -8.23 -16.24
g n	KIN Sha Obi Tot	.00 .00	.00 .03 .00 .00	. 82 . 85 . 88	- 20 - 20	. 02 . 00	.82 .20	69. 99. 99.	. 66	59. 95. 50.	. 22 . 22 . 23 . 23	99. 99. 99.	. 82 . 82 . 22	. 66 . 66 . 69
n h	KIN SNA OBI TOT	.10	.02 1.36 .00 1.36	1.46 .00 1.46	.00 .00	. 00 . 90	. 80 . 20	.90 -1.53 -1.57 -1.67	.00	.29 .90 99.	-1.53 15 -1.67	.82 -1.42 -,15 -1,57	1.36 .22 1.36	.20 26 15 21
ñh	kin Sna Oei Tot	.10	.00 1.36 .00 1.36	1.46 .00 1.46	. 89 . 36	.02 .02	. 00 . 00	-1.53 15 -1.67	.00	.22 .20 .00	-1.53 15 -1.67	.20 -1.42 15 -1.57	1.36 .22 1.36	.00 06 15 21
Bond Total	KIN SNA OBI TOT	-1.18 -1.18	-1.07 2.63 -3.34 -1.77	.39 -3.34 -2.95	-16.28 -16.23	11.36	-4.93 -4.93	-13.64 6.87 27 -7.60	.23	,2005 .12	-13.64 6.42 27 -7.48	-13.64 5.13 -16.55 -25.26	1.69 5.22 9.71	-13.64 6.81 -8.53 -15.36

Figure 15 (Continued)

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Figure 16. Binding energy decomposition and description of electron distribution for NH', BMAO approximation

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		Promotion		Qu clas	Quasi- classical		ring ration	Sharing Interference		
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
	KIN	- 4.1ó	-25.56		35.04					5.32
N ¹	NA	16.61	50.96		- 36. 32					31.25
	OEI	5.48	-25.12		3 70	11.66	-14 53			~12.80
	тот	17.94	. 28		2.42	11.66	-14.53			17.75
	KIN	. 00	13.43		-16.72					- 3.29
••	SNA	. 00	-11.14		8.90					- 2.24
н	OEI	. 00	. 00		4.59	14.34	- 3.23			15.69
	TOT	. 00	2.28		- 3.24	14.34	- 3 23			10.15
	KIN							-14.29		-14.29
	SNA			-1.07	3.13			6.21	. 11	8.37
BOND	OEI				- 5.42	-16.95	12.58	11		- 9,90
	тот			-1.07	- 2 29	-16,95	12.58	- 8.20	. 11	-15.82
	KIN	- 4 16	-12.13		18.32			-14.29		-12,26
	SNA	16.61	32, 82	-1.07	-24 29			6.21	. 11	37.38
NH	OEI	5.48	-25.12		2.87	2.05	- 5.18	11	-	-13.01
	TOT	17.94	2.56	-1.07	- 3.11	9.05	- 5.18	- 8.20	. 11	12.08

DENSITY PARTITIONING FOR VALENCE AO'S

	Valence Atomic Orbitals			Overlap	Bondorders and Populations							
	ĸ	s	σ	S(N. H)	ę	v	q	q ^N	q^T	^p (N, H)		
i -	999761	- 021121	005638	.037487	2.002	001	2.001	1. 998	. 003	030		
5	. 021798	. 982665	. 184102	. 505593	1.387	. 231	1.618	1.002	. 616	. 457		
e	.001652	164160	. 982891	. 312673	. 000	. 000	. 000	. 000	. 000	. 000		
Ħ		(unhybridize	d)	0	2.000	. 000	2.000	2.000	. 000	. 000		
h		(unhybridize	d)	-	. 151	. 230	. 381	1.000	619	-		

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

			N						
		i	ь	Ł	fr:	h			
	PR	2.002	004	. 000	. 000	-			
	SPN	. 000	.004	. 000	. 000	004			
1	SPT	. 000	. 000	. 00 0	. 000	. 00 3			
	TOT	2.001	.000	.000	. 000	001			
	PR	004	1.006	. 000	. 000	•			
t.	SPN	. 004	-1.204	.000	. 000	1.200			
D	SPT	. 000	1.507	. 000	. 000	891			
	тот	. 000	1.309	. 000	. 000	、309			
	PR	. 000	. 000	. 000	. 000	-			
	SPN	. 000	. 000	. 000	. 000	. 000			
1	SPT	. 000	. 000	. 000	. 000	. 000			
	TOT	. 000	. 000	.000	. 000	.000			
	PR	. 000	. 000	. 000	2.000	-			
	SPN	. 000	. 000	. 000	. 000	. 000			
1	SPT	. 000	. 000	. 000	. 000	. 000			
	TOT	. 000	. 000	. 000	2.000	.000			
	PR	<u> </u>	-	- ~ .	- ·	1.000			
L	SPN	004	1.200	. 000	. 000	-1.197			
п	SPT	. 003	891	. 000	. 000	. 270			
	TOT	001	. 309	. 000	. 000	. 073			

Figure 16 (Continued)

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	I						Prosot	on							Quasi-	Shu	ring	SP+QC
Orbital	PRH	k PRC	PRH	PRC	PRH O	PBC	PRH	TT PRC	рвн	F7 PRC	PBH	Total PEC	PB(H+C)	Orbital	classical QCT	Penet SPN	tration SPT	Total
KIN NA N k OEI TOT	-1.67 3.51 28 1.56	1.85 ~1.94 .21 .12	3.30 -3.46 -26.47 -26.64	.05 86 .25 .25	.00 .00 -24.39 -24.89	.00 .00 .08 .98	.00 .00 25.70 25.70	.00 .00 -3.84 -3.84	.00 .00 25.70 25.70	.00 .00 -3.84 -3.84	-1.65 3.49 63 1.20	1.89 -1.97 .39 .31	.23 1.52 24 1.51	KIN N 1 SNA N 1 OEI TOT	1.81 -2.74 .02 91	. 86	.01	1.81 -2.74 .89 84
KIN NA NOEI TOT	3.30 -3.46 -26.47 -26.64	.85 06 .25 .25	-68.77 208.77 -17.95 130.06	1.69 -2.52 .47 75	.00 .00 -17.78 -17.70	.00 .03 - 00 - 88	.00 .00 92 92	.00 .00 59 59	.88 .88 92 92	.00 .00 59 59	-54.28 201.97 -171.50 -23.93	1.77 -2.60 1.00 .17	-52.43 199.27 -178.51 -23.66	N D OEI TOT	33,23 -33,53 3,68 3,32	11.60 11.60	-14.54 -14.54	33,23 -33,58 .75 .39
KIN N OF OFI TOT	.00 .00 -24.89 -24.89	.00 .00 .08 .08	.09 .90 -17.79 -17.70	88. 98. 98 99	-48.29 176.32 1.67 129.78	.33 58 88 25	.00 .02 -3.22 -8.22	00. 00. 99 00	.00 .00 -3.22 -8.22	50. 50. 50 50	-48,29 176,32 -125,93 3,01	• 33 • • 58 • 12 • • 13	-47.95 175.74 -124.91 2.85	N L SNA N L OEI TOT	. 00 . 00 . 00 . 00	08 00	.00 .88	. 20 . 20 28 92
KIN N TOEI TOT	.00 .00 25.70 25.70	.00 .00 -3.84 -3.84	.00 92 92	.00 .00 59 59	.00 .00 -8.22 -5.22	.00 .00 00 00	49,99 -182,53 19,92 -112,63	-14.77 28.05 -1.41 11.86	.99 .90 27.94 27.94	.00 .00 -2,76 -2,76	49.99 -182.53 151.32 18.78	-14.77 28.05 -13.31 -03	35.21 -154.48 138.81 18.74	H TT OEI TOT	. 80 . 99 . 99 . 99	00 00	. 86 . 88	. 88 . 99 80 69
KIN N TO TOT	.00 25.78 25.78	.00 .00 -3.84 -3.84	.09 .72 92	.60 .88 59 59	.00 .89 -8.22 -8.22	.00 .00 00 00	.00 .00 27.04 27.04	.00 -2.76 -2.76	49.99 -182.53 19.92 -112.63	-14.77 28.05 -1.41 11.86	49.99 -192,53 151.32 18.78	-14.77 28.05 -13.31 03	35.21 -154.48 138.81 	KIN N 77 SHA DEI TOT	. P0 . A0 . DD . DD	60 08	.00	. 80 . 80 - , 80 - , 68
Total N KIN GEI TOT											-4,16 16,61 5,48 17,94	-25.56 50.96 -25.12 .28	-29,72 67,57 -19,63 13,21	Total N KIN SNA OEI TOT	35.84 -36,32 3.70 2,42	11.66 11.66	-14.53 -14.53	35,84 -36,32 ,83 -,46
Total H KIN 051 TOT											. 60 . 66 . 66 . 66	13.43 -11.14 .00 2.28	13.43 -11.14 .00 2.28	Total H KIN SNA OEI TOT	-16,72 8,90 4,59 -3,24	14.34 14.34	-3.23 -3,23	-16.72 8.90 15.69 7.87

INTRA-ATONIC ENERGY CONTRIBUTIONS OF NR*, BHAU

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INTERATORIC ENERGY CONTRIBUTIONS OF MRT, BRAD

Orbita	a 1	Quas	iclassical Electro	static				aring				1	- Totals -	
Pain M H	r	QCN	QCT	QC(H+T)	. SPN	enetration SPT	SP(N+T)	SIN		- Interference	SI(N+T)	И	T	N+T
1 h	KIN SNA OBI TOT	e3 03	00 .02 02 00	81 82 83	.05 .85	84 84	. 01 . 01	.06 02 .77	62	0101 ø2	.73 .82 89 .75	.73 .02 .85 .80	.00 07 07	.73 .02 02 .73
ъ в	KIN SNA OEI TOT	-,88 -,88	<i>1.6</i> 37 -5.48 -6.43	-1.91 -5.4P -7.30	-17.01 -17.01	12.63 12.63	-4.38 -4.38	+15.83 7.29 -6.73	•97	.1905 .14	-15.83 8.39 .93 -6.60	-15.03 7.38 -16.97 -24.62	98 7.23 6.33	-15.83 6.49 -7.74 -18.28
£ h	KIN SHA OBI TOT	.0P .50	.00 .00 .60 .50	. 88 . 110 . 89	.ev .ev	. ee . ov	. 80 . 82	.00 .08 .00 .00 .00	.ee	.00 .60	.80 .90 .90 .99	.08 .00 .00 .00	. 88 . 66 . 60	. 110 . 68 . 69 . 69
πh	KIN SNA CEI TOT	•.09 09	.00 2.07 .00 2.07	1.98 ,00 1.98	. re . ee	. 88 . 80	. eb . eb	.05 -1.05 07 -1.12	.80	.00 .00 . 0 9	.00 -1.85 67 -1.12	.00 -1.14 07 -1.21	2.87 .00 2.47	.00 .93 07 .55
ñħ	KIN SNA OEI TOT	89 89	.00 2.07 .00 2.07	1.98 .00 1.98	. Bit . Bit	. 2C . DN	. ee . 08	.00 -1.05 07 -1.12	.80	.00 .00 .00	.00 -1.03 07 -1.12	.00 -1.14 07 -1.21	2.07 .00 2.07	.88 .73 87 .86
Bond Total	KIN SNA OEI TOT	-1.07	-,66 3,79 -5,42 -2,29	2.05 -5.42 -3.37	-16.95 -16.95	12.58 12.58	-4.37 -4.37	-14.29 5.26 11 -3.28	.95	.1786	-14.29 6.32 11 -8.09	-14.29 5.13 -17.86 -26.23	3.24 7.17 18.41	-14.29 8.37 -7.98 -15.82

Figure 17. Binding energy decomposition and description of electron distribution for NH", SAO approximation

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BINDING ENERGY PARTITIONING

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		Promotion		Quasi- classical		Sharing Penetration		Sharing Interference			
_		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL	
N	KIN NA OEI TOT	- 5.49 17.44 7.63 19.57	- - -		27.95 -27.47 2.16 2.64	12.19 12.19	-12.82 -12.82			22.46 -10.03 9.15 21.57	
н	KIN SNA OEI TOT		- - -		- 6.43 4.83 1.90 .29	10.64 1064	- 3.33 - 3.33			- 6.43 4.83 9.22 7.61	
BOND	KIN SNA OEI TOT			-1.03 -1.03	1.23 - 2.78 - 1 55	-15.58 -15.58	10.76 10.76	-14.95 7.12 24 - 8.07	. 15 . 15	-14.95 7.47 - 7.85 -15.32	
NH	KIN SNA OEI TOT	- 5 49 17 44 7 63 19 57	- - -	-1.03 -1.03	21.52 -21.41 1.28 1.39	7.25 7.25	- 5.39 - 5.39	-14.95 7 12 24 - 8.07	. 15 . 15	1.08 2.27 10.52 13.86	

DENSITY PARTITIONING FOR VALENCE AO'S

	Valence Atomic Orbitals			Overlap	Bondorders and Populations						
	k	S	σ	S _(N, H)	p	v	٩	q ^N	q ^T	^р (N, H)	
i	. 299855	016544	004081	. 045305	2.003	002	2.001	1. 298	. 003	042	
ь	. 016988	, 366442	. 163227	. 534668	1.186	. 280	1.472	1.002	. 470	. 535	
l	. 001 325	163273	, 986580	292164	. 000	. 000	. 000	. 000	. 000	. 000	
π		(unhybridize	d)	0	2.000	. 000	2.000	2.000	.000	. 000	
h	1	(unhybridize)	d)	-	. 243	. 284	. 527	1.000	473	-	

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

			N				
		i	ь	Ł	π	h	
	PR	2.002	004	. 000	. 000	-	
	SPN	. 000	. 00 5	. 000	. 000	005	
1	SPT	. 000	. 000	. 000	. 000	. 003	
	тот	2.002	.000	.000	. 000	001	
	PR	004	1.007	. 000	. 000	-	
	SPN	. 005	-1.261	. 000	. 000	1.256	
Ь	SPT	. 000	1.338	. 000	. 000	868	
	тот	. 000	1.084	. 000	. 000	. 389	
	PR	. 000	. 000	. 000	. 000		
	SPN	. 000	. 000	. 000	.000	. 000	
X	SPT	. 000	. 000	. 000	. 000	. 000	
	тот	. 000	. 000	. 000	.000	. 000	
	PR	. 000	. 000	. 000	2.000	-	
	SPN	. 000	.000	.000	. 000	. 000	
17	SPT	. 000	. 000	. 000	. 000	. 000	
	тот	. 000	. 000	.000	2.000	. 000	
	PR	-	-	-	-	1.000	
	SPN	005	1.256	. 000	. 000	-1.252	
h	SPT	. 003	868	. 000	.000	. 391	
	TOT	001	. 389	. 000	. 000	. 139	

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INTRA-ATOMIC ENERGY CONTRIBUTIONS OF NH*, SAO

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	I						1	Quasi-	Sh	aring	SP+QC
Orbital	k		σ	fromotion fr	ŤŤ	Total	Orbital	QCT	SPN	SPT	Total
KIN N k OEI TOT	-1,44 3.01 26 1,31	2.65 -2.77 -26.61 -26.73	.22 .02 -25.50 -25.52	.02 .20 26.15 26.15	.00 .00 26.15 26.15	-1.43 3.82 57 1.00	N 1 OEI TOT	1.90 -2.85 .02 92	.10	.01	1.92 -2.85 .13 81
KIN N S OEI Tot	2.65 -2.77 -26.61 -26.73	-62.45 218.60 -17.73 132.41	.80 .80 -17.96 -17.96	55. 58. 12 13	.22 .02 51 51	-57.17 205.08 -157.90 -3.99	KIN N b OSI TOT	26.05 -24.62 2.14 3.56	12.89 12.89	-12.83	26.25 -24.62 1.39 2.82
KIN N O OEI TOT	.ee .ee -25.5e -25.5e	.20 .02 -17.96 -17.96	-52.34 180.71 1.71 132.88	.02 .20 -8,52 -8,52	.02 .02 -3.50 -6.52	-58.34 182.71 -123.64 2,33	SNA SVA SVA OEI TOT	.02 .00 .02 .02	00 00	. <i>30</i> .00	.00 .20 22 00
KIN N T CEI TOT	.20 .00 26.15 26.15	03. 50. 13 13	.00 .00 -8.50 -8.50	51.72 -135.68 20.45 -113.50	.20 .20 27.51 27.51	51.72 -185.68 147.27 13.12	N T OEI TOT	60. 50. 99.	22 22	. 22 . 63	. 23 . 99 92 28
KIN N 77 OEI TOT	.00 26.15 26.15	.23 .90 61 61	.C2 .00 -8.52 -8.50	.92 .02 27.51 27.51	51.72 -185.63 20.45 -113.52	51.72 -165.65 147.07 13.12	B T SNA DEI TOT	. 80 . 80 . 00 . 00	02 00	.22. 59.	. ee . ee ee ee ee
Total N KIN NA OEI TOT						-5.49 17.44 7.63 19.57	Total N KIN SNA OEI TOT	27.95 -27.47 2.16 2.64	12.19	-12,92 -12,92	27.95 -27.47 1.52 2.82
							Total H KIN SWA OEI TOT	-5.43 4.83 1.92 .29	10.64	-3.23 -3.33	-6.43 4.63 9.22 7.61

INTERATORIC	ENERCY	CONTRIBUTIONS	QP	₩₩°,	540	
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Orbital Pair	Qua	siclassical Electro	ostatic	Sharing					••••••••••••••••••••••••••••••••••••••		- Totals			
Pai N H	r	QCH	्रता	QC(N+T)	SPN	enstratio SPT	SF(N+T)	SIN		- Interference SIT	SI(E+T)	×	T	N+T
i h	KIN SNA OEI TOT	03 03	00 .02 02 00	22 22 23	.26	24 24	.21 .21	23 23 21 1.13	22	6182 61	1.15 07 01 1.11	1.19 29 .05 1.15	26 26 26	1.19 29 01 1.29
bh	KIN Sha Oei Tot	99 99	7312 -2.77 -3.61	-1.34 -2.77 -4.61	-15.64 -15.64	12.32 10.30	-4.84 -4.84	-16.14 9.63 -6.00	.47	.22 •.25 .17	-16.14 12.27 .24 -5.83	-16.14 9.11 -15.60 -22.63	68 8.03 7.35	-16.14 8.43 -7.56 -15.27
lh	KIN Sha Oei Tot	.ee .øø	.00 .00 .00	. 88 . 88 . 20	.20 .00	. 22 . 22	.ee	98. 29. 29.	. 20	.ce .ee .øø.	. 29 . 20 . 20 . 22	.88 .88 .82 .82	. 00 . 22 . P e	.22 .22 .22 .00
πħ	KIN SNA OEI TUT	.00 .00	.20 1.23 .22 1.23	1.03 .20 1.03	. 20 . 20	. ee . ee	. 22 . 20	.00 -1.47 14 -1.60	.20	.99. 99.	-1.47 -14 -1.62	.28 -1.47 14 -1.62	1.23	.00 43 14 57
₩þ	KIN SNA OEI TOT	.00 .00	.82 1.83 .88 1.83	1.93 .22 1.23	. 22 . 22	. 22 . 22	.22.	.23 1.47 -1.4 -1,62	.82	.22 .22 .22	-1.47 -1.47 -1.66	.22 -1.47 14 -1.62	1.23 .28 1.23	- 43 - 14 - 57
Bond Total	KIN SNA OEI TOT	-1.03	73 1.96 -2.78 -1.55	.21 -2.78 -2.58	-15.58	12.76	-4.82 -4.82	-14.95 6.57 24 -8.07	.45	.21 •.26 .15	-14.95 7.27 24 -7.92	-14,95 6,79 -15.83 -24,68	1.35 7.97 9.36	-14.95 7.47 -7.95 -15.32



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		Promo	otion	Qua	asi- sical	Sha: Penet	ring ration	Shari Interfe	ing rence	
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
N	KIN NA OEI TOT	- 4.20 16.48 6.78 19.06	- - -		28.69 -28.63 2.52 2.58	11.85 11.85	-13.09 -13.09			24.49 -12.15 8.06 20.40
н	KIN SNA OEI TOT		-		- 7.01 5.25 2.25 .50	10.50 10.50	- 3,00 - 3.00			- 7.01 5.25 9.75 8.00
BOND	KIN SNA OEI TOT			-1.04 -1.04	1.40 - 3.30 - 1.90	-15.37 -15.37	10.80 10.80	-13.69 6.33 21 - 7.57	. 15 . 15	-13,69 6.85 - 8.08 -14.92
NH	KIN SNA OEI TOT	- 4.20 16.48 6.78 19.06	-	-1.04 -1 04	21 68 -21.98 1 47 1 17	6.98 6.98	- 5.29 - 5.29	-13.69 6.33 21 - 7.57	. 15 . 15	3.79 - 05 9.73 13.47

DENSITY PARTITIONING FOR VALENCE AO'S

	Valence Atomic Orbitals			Overlap Bondorders and Populations						
	k	5	σ	^S (N, H)	р	v	q	qN	γT	P(N, H)
i	. 000784	020294	004537	. 043615	2.003	002	2.001	1. 998	. 00 3	038
ь	. 020765	.986065	. 165060	. 542155	1.238	. 276	1.514	1.002	512	. 510
L	.001124	165118	. 986273	. 298061	. 000	. 000	. 000	. 000	. 000	. 000
11		(unhybridize	d)	0	2.000	. 000	2.000	2.000	.000	. 000
h	l	(unhybridize)	d)	-	, 210	. 275	. 485	1.000	515	-

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

		I		н		
		i	ь	£	π	h
	PR	2.002	004	. 000	. 000	-
	SPN	. 000	.004	. 000	. 000	004
1	SPT	.000	. 000	. 000	. 000	003
	тот	2.002	. 000	.000	. 000	001
	PR	004	1.006	. 000	. 000	-
1	SPN	. 004	-1.244	.000	. 000	1.239
D	SPT	. 000	1.384	. 000	. 000	871
	тот	. 000	1.146	.000	. 000	. 368
	PR	. 000	. 000	.000	. 000	-
^	SPN	.000	. 000	.000	. 000	. 000
x	SPT	. 000	. 000	.000	. 000	. 000
	TOT	.000	. 000	.000	. 00'-	.000
	PR	. 000	. 000	.000	2.000	-
-	SPN	. 000	. 000	.000	. 000	. 000
77	SPT	.000	.000	.000	.000	.000
	тот	. 000	. 000	.000	2.000	. 000
	PR	-	•	-	-	1.000
	SPN	004	1.239	.000	. 000	-1.235
n	SPT	. 003	871	. 000	.000	. 353
	TOT	001	. 368	.000	. 000	. 118

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							1	Quasi-	Sharing		SP+QC
Orbital	k		or	from ton	fr	Total	Orbital	QCT	SPN	SPT	Total
N KIN N K OBI TOT	-1.54 3.23 25 1.44	3.19 -3.35 -26.26 -26.42	.20 .22 -25.26 -25.06	.22 .02 25.71 25.71	.22 .22 25.71 25.71	-1,52 3,21 -,53 1,12	N 1 CEI TOT	1.96 -2.90 .22 92	.22	. 21 . 21	1. E6 -2. 30 .11 53
KIN N B OEI TOT	3.19 -3.35 -26.26 -26.42	-60.39 227.46 -17.62 129.45	.20 .02 -17.70 -17.70	.78 .22 56 66	.02 .52 56 66	-54,23 222,79 -161,69 -14,93	KIB SNA N D OEI TOT	26.93 -25.83 2.50 3.52	11.77	-13,12 -13,12	26.83 -25.33 1.17 2.17
KIN N G OEI TOT	.00 .00 -25.06 -25.26	.20 .20 -17.70 -17.70	-48.62 177.54 1.65 130.60	.22 .22 -8.34 -8.14	.22 .02 -8,34 -8,34	-48,62 177,54 -125,87 3,25	N CEI TCT	.00 .02 .02 .00	22 00	. 22 . 20	55. 55. - 20 00
KIN N TT OEI TOT	.80 .02 25.71 25.71	.22 .22 66 66	.22 .02 -8.34 -8.34	49.99 -182.53 20.26 -1:2.48	.20 .22 27.24 27.24	49,59 -182,53 147,46 14,91	N T OEI TOT	.22 .99 .20 .20	22 66	.ec .ce	59. 55. 25 25
KIN NFOEI TOT	.00 .00 25.71 25.71	55. 29. 23 23 23	.22 .02 -8.34 -8.34	.22 .20 27.04 27.04	43.99 -182.53 27.26 -112.48	45,99 -152,53 147,45 14,91	N T SNA N T SNA OEI TOT	.00 .02 .22 .90	00 00	. ee . ee	. 22 . 22 22 20
Total N KIN NA CEI TOT						-4.20 16.48 6.75 19.26	Total N KIN SNA GEI TOT	28.69 -28.63 2.52 2.58	11,35	-13.29 -13.29	26.69 -28.63 1.28 1.34
							Total H KIE SNA QEI TOT	-7.01 5.25 2.25 .50	10.50	-3.20 -3.22	-7.21 5.25 9.75 8.22

INTRA-ATOMIC ENERGY CONTRIBUTIONS OF SH*, BAO

INTERATORIC	ENERGY	CONTRIBUTIONS	٥P	5H*,	Báú

Orbital		Quasiclassical Electrostatic			Sharing								Totals		
Pai N H	F	QCN	ącr	20(N+T)	SPN	Penetratic SPT	SP(N+T)	SIN	Interference SIT	SI(N+T)	N	7	N+T		
ih	KIN SNA OEI TOT	03 23	00 02 03	22 02 23	. 25 . 85	24 24	. 21 . 21	.25 21 1.95	22212P -,81	1.73 .82 01 1.24	1.23 02 .05 1.28	- 22 - 26 - 26	1.23 22 21 1.22		
bh	KIN SNA OEI TOT	-1.21	5113 -3.28 -4.21	-1.95 -3.28 -5.23	-15.42	12.94 12.94	-4.58 -4.58	-14.72 E.66 .03 -5.59	44 .2226 .16	-14.72 9.27 .23 -5.42	-14.72 8.09 -15.39 -22.02	75 7.56 6.82	-14.72 7.32 -7.52 -15.23		
<i>k</i> h	KIN SNA OEI TOT	.0e .00	.02 .09 .09 .09	.22 .20 .33	. 20 . 22	.27 .22	.20 .20	99. 99. 99.	95. 52. 55 68 .	.00 .22 .20 .20	.02 .22 .20 .22	. P2 . 32 . 72	-22 -22 -42 -22		
nk	KIN SNA OEI TOT	.01	.00 1.16 .02 1.16	1.16 .22 1.16	.ce .ce	. 20 . 20	. 22 . 22	-1.40 .0 12 -1.52	95. 92. 59 88.	.22 -1.40 12 -1.52	-1.39 -1.2 -1.51	1.15 .22 1.16	.22 23 12 36		
ˈ77 h	KIN SNA OEI TOT	.e1 .01	.00 1.16 .00 1.16	1.16 .22 1.16	. 89 . 82	. 22 . 22	. 00 . 20	.00 -1.42 .0 -1.52	96 . 53. 55	-:.42 -:.42 12 -1.52	.20 -1.39 12 -1.51	1.16 .02 1.16	.22 23 12 36		
Bond Total	KIN SNA GEI TOT	-1.84 -1.24	81 2.21 -3.32 -1.90	.37 -3.30 -2.93	-15.37	12.82	-4.57	-13.69 5.91 .4 21 -7.57	•2 •21 •.26 •15	-13.65 6.49 21 -7.42	-11.69 5.12 -15.57 -23.5	:.55 7.52 9.26	-13.69 6.35 -3.23 -14,92		



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Figure 19. Binding energy decomposition and description of electron distribution for NH", BMAO approximation

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		Promotion		Quasi- classical		Sha Penet	Sharing Penetration		ng rence		
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL	
N	KIN NA OEI TOT	- 4.71 16.98 4.94 17.21	-28.69 56.58 -27.48 .42		36.26 -37.75 4.03 2.54	11.40 11.40	-14.74 -14.74			2.87 35.82 -21.84 16.84	
н	KIN SNA OEI TOT	.00 .00 .00 .00	11, 46 - 9, 72 .00 1, 74		-16.29 9.00 4.87 - 2.42	13.70 13.70	- 2.87 - 2.87			- 4.83 72 15.70 10.16	
BOND	KIN SNA OE1 TOT			88 88	2.68 - 5.56 - 2.88	-15.63 -15.63	11.83 11.83	-13.15 5.87 08 - 7.36	. 14 . 14	-13.15 7.80 - 9.44 -14.78	
NH	KIN SNA OEI TOT	- 4.71 16.98 4.94 17.21	-17.23 46.86 -27.48 27.15	88 88	19.97 -26.07 3.34 - 2.76	9.47 9.47	- 5.78 - 5.78	-13.15 5.87 08 - 7.36	. 14 . 14	-15.11 42.90 -15.58 12 22	

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DENSITY PARTITIONING FOR VALENCE AO'S

	Valence Atomic Orbitals			Overlap		Bond	orders and	and Populations			
	k	s	σ	^S (N, H)	Р	v	q	q ^N	\mathbf{T}_{p}	^p (N, H)	
i	. 999816	018639	004580	. 0 32278	2.001	001	2.001	1.998	. 003	026	
ь	.019140	.986018	. 165536	. 462178	1.440	. 209	1.649	1.002	. 647	. 452	
l	. 001431	165593	. 986193	. 308842	. 000	. 000	. 000	. 000	. 000	. 000	
77	(unhybridized)			0	2.000	. 000	2.000	2.000	. 000	. 000	
h		(unhybridize)	d)	-	. 142	. 208	. 350	1.000	650	-	

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

			N						
		i	ь	L	π	h			
	PR	2.001	003	. 000	. 000				
i	SPN	. 000	. 003	. 000	. 000	003			
•	SPT	. 000	. 000	. 000	. 000	. 00 3			
	TOT	2.001	. 000	. 000	. 000	001			
	PR	003	1.005	. 000	. 000	-			
L.	SPN	. 003	-1.194	. 000	. 000	1.190			
D	SPT	. 000	1.548	. 000	. 000	901			
	тот	. 000	1.360	. 000	. 000	. 289			
	PR	. 000	. 000	. 000	. 000	-			
	SPN	. 000	. 000	.000	. 000	. 000			
X	SPT	. 000	. 000	. 000	. 000	. 000			
	тот	. 000	. 000	. 000	. 000	. 000			
	ਸ਼ਕ	000	. 000	. 000	2,000				
	SPN	. 000	. 000	. 000	. 000	, 000			
r	SPT	. 000	. 000	. 000	. 000	. 000			
	TOT	. 000	. 000	. 000	2.000	. 000			
	ब्रद	_	-	_	-	1,000			
	SPN	- 003	1,190	. 000	. 000	-1.187			
h	SPT	003	- 901	000	. 000	. 249			
	TOT	- 001	289	000	. 000	. 062			
	101	= .001	. 20 /						
Figure 19 (Continued)

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INTEA-ATONIC ENERGY CONTRIBUTIONS OF NH", BRAD

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1	Prosotion								ſ	Quasi-	Shu	ring	SP+QC					
Grbital	PBH	PAC	PRH	PRC	PRH	PRC	PRH	PRC	PBH	PRG	PBR	PRC	PB(R+C)	Orbital	QCT	SPM	SPT	Total
KIN N R NA OEI TOT	-1,50 3,15 -,25 1,40	1.99 -2.89 .24 .14	2,91 -3,86 -26,35 -26,49	.82 82 .89 .09	.00 -25.86 -25.86	.88 .80 .84 .84	.00 .00 25.71 25.71	.08 .00 -4.10 -4.13	.00 .00 25.71 25.71	.00 .00 -4.18 -4.10	-1,49 3,13 -,57 1,88	2.82 -2.11 .42 .33	.53 1.03 14 1.41	N 1 OBI TOT	1.60 -2.42 .82 80	.05	.00	1.60 -2.42 .68 75
KIN NA NBOBI TOT	2.91 -3.86 -26.35 -26.49	.#2 #2 .#9 .#9	-68.48 207,51 -18.09 129.02	.56 85 .10 19	.80 .08 -17.78 -17.78	.88 .88 88 00	.00 .00 97 -,87	.00 .08 69 69	.00 .00 87 87	.00 .88 69 69	-54.59 201.40 -186.68 -39.86	.57 36 .39 .10	-54.01 200.54 -136.28 -39.76	KIN N b SNA DEI TOT	34.66 -35.32 4.01 3.35	11.35 11.35	-14.74 -14.74	34.66 -35.32 .62 84
KIN N J OEI TOT	.00 .00 -25.06 -25.06	. 60 . 80 . 84 . 84	.00 .00 -17,70 -17,70	68. 98. 98 98	-48.61 177.51 1.67 138.57	.16 28 00 12	.00 .00 -8,34 -8,34	.00 .00 01 01	.80 .90 -3.34 -8.34	.00. .00 01 01	-48.61 177.51 -125.87 3.03	.16 28 .14 .02	-48.46 177.23 -125.73 3.85	N & SNA N & OEI TOT	. 98 . 89 . 99 . 89	80 88	.80 .80	. 88 . 88 88 88
KIN NA N TOEI TOT	.00 .00 25,71 25,71	.00 .70 -4.10 -4.10	.80 .80 87 87	.08 .00 69 69	.00 .00 -8,34 -8,34	.00 .00 01 01	49.99 -182.53 19.86 -112.69	-15.71 29.92 -1.50 12.70	.80 .90 27.84 27.84	.88 .88 -2.94 -2.94	49,99 -182,53 159,02 26,48	-15.71 29.92 -14.22 02	34.27 -152.62 144.80 26.46	N TY OEI TOT	.00 .92 .60 .00	88 88	.80 .80	. 66 . 66 66 66
KIN NA NA OR1 TOT	.88 .98 P5.71 25,71	.08 .08 -4.10 -4.10	.00 .08 87 87	.00 .00 69 69	.08 .08 -8.34 -8.34	.00 .00 .01 01	.00 .09 17.04 27.04	.00 .00 -2.94 -2.94	49.99 -192.53 19.86 -112.69	-15.71 29.92 -1.50 12.70	49.99 -182.53 159.82 26.48	-15.71 29.92 -14.22 02	34.27 -152.62 144.88 26.46	KIN N 77 OBI TOT	.98 .99 .99 .99	08 08	.00	.08 .08 08 08
Total N XIN NA ORI TOT											-4.71 16.98 4.94 17.21	-28.69 56.58 -27.48 .42	-33,39 73,57 -22,54 17,63	Total N KIN SNA OEI TOT	36.26 -37.75 4.83 2.54	11.48 11.40	-14.74 -14.74	35.26 -37.75 .78 79
Total H KIN NA OBI TQT											. 68 . 89 . 69 . 60	11.46 -9.72 .00 1.74	11.46 -9.72 .08 1.74	Total H KIN SNA OEI TOT	-16.29 9.80 4.87 -2,42	13.70	-2.87 -2.87	-16.29 9.20 15.70 8.42

INTERATONIC ENERGY CONTRIBUTIONS OF NH*, BRAG

Ort	ite	a	Qua	siclassical Electro	ostatic			- ··· St	aring					- Totals -	
_ »	alr H		QCN	QCT	QC(N+T)	SPN	SPT	SP(N+T)	SIN		SIT	SI(N+T)	ы	T	X+T
1	ħ	KIN SNA OBI TOT	02 02	00 .02 02 01	00 02 82	.04 .24	84 C4	.e1 .01	.84 00 .57	01	8181 82	.55 .01 ~.08 .56	.55 .00 .24 .60	06 86	.55 .89 82 .54
ъ	h	KIN SNA OEI TOT	68 68	- .54 36 -5.54 -6.43	-1.58 -5.54 -7.11	-15.67 -15.67	11.8C 11.86	- 5. 81 - 3. 81	-13.71 6.75 .03 -5.97	.95	.1802 ,16	-13,71 7,85 .03 -5,82	-13,71 7,02 -15,64 -22,32	75 6.33 5.58	-13.71 6.28 -3.51 -16.74
R	h	KIN SNA OE1 TOT	.00 .80	.00 .00 .90 .00	.00 .00 .00	.02 .09	. 11H . CB	. 64 . PC	.48 .84 .80 .80	. 88	.en .op .bn	.00 .00 .00	.80 .00 .00 .00	.00 .00 .00	.88 .83 .83 .88
•	h	KIN SNA OEI TOT	89 89	.00 1.78 .00 1.78	1.69 .00 1.69	.ee .ce	. ve . ee	.00 .00	.93 93 95 98	.00	an, An,	.118 93 85 98	.00 -1,02 05 -1,07	1.78 .00 1.78	.00 .76 05 .71
,	h	KIN SNA OEI TOT	89 89	.00 1.79 .00 1.78).69 .0P 1.69	.00 .00	. 60 . Fu	. 88 . 88	.93 93 05 98	.00	.00 .98 .85	.93 93 95 98	.00 -1.02 25 -1.07	1.78 .00 1.78	.00 .76 05 .71
В	ond tal	KIN SNA OEI TOT	• .88 88	54 3.22 -5.56 -2.88	1.80 -5.56 -3.75	-15.63	11.83	-3.80	-13.15 4.93 08 -7.36	.44	.1783 .14	-13.15 6.04 08 -7.23	-13.15 4.99 -15.71 -23.87	2.82 6.27 9.89	-13.15 7.80 -9.44 -14.78

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Figure 20. Binding energy decomposition and description of electron distribution for HF, SAO approximation

BINDING ENERGY PARTITIONING

		i Prom	otion	Qua	asi- sical	Sha Pene	ring tration	Shari Interíe	ing rence	
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL
F	KIN NA OEI TOT	1,88 ,45 - ,11 2,22	- - -		11.90 -10.70 31 1.51	7.01 7.01	-1.78 -1.78			13.78 -10.25 5.44 8.97
н	KIN SNA OE1 TOT	-	- - -		- 2.07 1.57 .20 32	3.33 3.33	1.52 1.52			- 2.09 1.57 5.06 4.53
BOND	KIN 5NA OEI TOT			-1,72 -1,72	. 68 34 33	-ú 14 -0.14	- 61 61	-11.39 6.19 -1.14 -0.34	29 - 29	-11, 39 4, 84 - 8, 23 -14, 78
HF	KIN SNA OEI TOT	1 88 . 45 11 2.22	- - -	-1,72 -1,72	9,81 - 8,45 ,17 1 52	4. 20 4. 20	- 87 67	-11 33 6.19 - 1.14 - 6.34	-,29 -,29	. 30 - 3.82 2.25 - 1.28

DENSITY PARTITIONING FOR VALENCE AO'S

	Vale	nce Atomie Or	rbitals	Overlap		Bond	orders and	ons		
	k	\$	5	^S (F, H)	p	v	ą	q ^N	q ^T	^p (F,H)
i .	, 291435	128805	. 021583	. 000000	2.000	. 000	2.000	2.000	. 000	. 000
l	. 128 347	. 939356	343400	. 343688	2.234	175	2.059	2.003	. 056	509
ь	. 024159	. 345288	. 938919	. 444080	. 789	. 306	1.095	5.92	. 0 ?8	. 688
tr		(unhybridize	d)	0	2.000	. 000	2.000	2.000	. 000	. 000
h.		(unhybridize	d)	•	716	. 131	. 846	1.000	154	-

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

				- F		н
		i	R	ь	₩	h
i	PR SPN SPT TOT	2.000 .000 .000 2.00 0	.000 .000 .000 .000	.000 .000 .000 .000	. 000 . 000 . 000 . 000	. 000 . 000 . 000
2	PR SPN SPT TOT	.000 .000 .000 .000	2.063 004 .100 2.159	060 .043 .004 013	.000 .000 .000 .000	039 048 087
b	PR SPN SPT TOT	.000 .000 .000 .000	060 .043 .004 013	1.057 474 .021 .604	.000 .000 .000 .000	. 431 . 073 . 504
fr	PR SPN SPT TOT	.000 .000 .000 .000	.000 .000 .000 .000	.000 .000 .000 .000	2.000 .000 .000 2.000	- .000 .000 .000
h	PR SPN SPT TOT	. 000 . 000 . 000	039 048 087	. 431 . 073 . 504	. 000 . 000 . 000	1.000 392 179 .429

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	{			reation			1	Quasi-	Sha	ring	SP+QC
Orbital	×	· 8	σ.	17	₩	Total	Orbital	QCT	SPN	SPT	Total
F k OEI TOT	•.55 1.14 •.17 •42	2.27 -2.35 -3.79 -3.87	.28 .03 4.89 4.69	. 22 . 22 22 22	. 28 . 29 22 22	53 1.12 32 .27	F 1 OEI TOT	.00 .00 .00	00 00	. 86	- 20 - 20 - 20
KIN Psoei Tot	2.27 -2.55 -3.79 -3.87	-12.79 40.97 -3.22 24.95	.20 .20 .14 .14	.20 .20 -3.21 -3.01	-3.21 -3.21	-8,27 36,29 -23,13 4,89	P L SNA P L OEI TOT	3,10 -3,47 ,11 -,27	.19	-1.42 -1.42	3.10 -3.47 -1.13 -1.50
F J DEI TOT	.00 .00 4.09 4.09	.00 .02 .14 .14	12.68 -36.96 3.16 -23.12	.20 .00 2.88 2.88	.00 .00 2.88 2.88	12.68 -36.96 23.49 -2.79	F b OSI TOT	8.90 -7.23 .28 1.78	6.83 6.83	35 35	8.80 -7.23 6.68 8.25
RIF P T OEI TOT	.80 .00 62 62	.20 .00 -3.01 -3.01	.20 .00 2.55 2.58	. 00 . 00 . 23 . 25	. 00 . 00 . 02 . 02	. 00 . 60 08 28	P W OEI TOT	.00 .00 .00	00 00	.00	.00 00 00
P = NA OEI TOT	.00 .00 02 02	.22 .22 -3.21 -3.21	.20 .00 2.35 2.58	55. 56. 56. 59.	. 80 . 00 28 26	. UC . 22 09 28	P T SNA DEI TOT	.00 .00 .00	00 08	. 90 . 90	99. 99. 93 35
Total F KIN NA QEI TOT			,,,,,			1.58 .45 11 2.22	Total P EIN SNA QEI TOT	11.90 -10.70 .31 1.51	7.01 7.81	-1.78 -1.78	11.90 -10.70 5.55 6.75
							Total H KIN SNA OEI TOT	-2.89 1.57 .20 32	3.33	1.52	-2.29 1.57 5.26 4.53

INTRA-ATOMIC ENERGY CONTRIBUTIONS OF HP, SAC

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INTERATORIC	ENERGY	CONTRIBUTIONS	OP	HP.	SAO

Orbita	1	Que	Quasiclassical Electrostatic									Totals			
Pain P H	r	QCN_		201		QC(N+T)	SPN	SPT	SP(N+T)	SIN	Interference	SI(N+T)	¥	Ť	H+T
1 h	KIN SNA OEI TOT	e5 e5	.02	. 22 . 21	.eı	84 .02 24	.02 .29	. 20 . 20	. 22	.12 .12 .12 .03	0 .20 .00 ,60	.00 .12 10 .03	.30 .27 10 82	. 21 . 22 . 21	.82 .08 12 01
f n	KIN SNA OEI Tot	1.12	24	10 .61	.75	1.85 - 12 1.73	.63 .63	• 57 • 57	1.22 1.22	10.75 -1.4078 12 8.46	1102 13	12.75 -2.31 12 8.32	10.75 -1.06 .51 10.20	.58 .47 1.05	10.75 48 .98 11.25
bh	KIN SNA OEI TOT	-3.13 -3.13	44	- 24 - 95	.27	-3.84 -,24 -4.06	-6.77 -6.77	-1.17 -1,17	-7.95 -7.95	-22.14 9.72 -2.43 .15 -14.72	.et25 - .16	-22.14 7.11 .15 -14.88	-22.14 4.14 -6.62 -24.62	87 -1.42 -2,28	-22.14 3.27 -8.83 -26.91
πh	KIN SNA OEI TOT	.17 .17	.02	.ee .33	.33	.52 .62 .50	.ee .ee	. ee . ee	. 80 . 82	.80 .49 .22 54 05	. 22 . 22 .00	.02 .49 54 25	.28 .66 54 .12	.33 .20 .33	. 20 . 98 54 . 44
₩h	KIN SNA OEI TOT	•17 •17	. 89	. 88 . 33	.33	. 5e . 20 . 5e	. 22 . 83	. 22 . 92	. 80 . 80	.80 .49 .00 54 05	.09 .09	.88 .49 54 85	.00 .66 54 .12	. 33 .00 .33	.20 .98 54 .44
Bond Total	KIN SNA OEI TOT	-1.72	48	34	, 16	-1,25 34 -1,39	-6.14 -6.14	61 61	-6.75 -6.75	-11.39 9.42 -3.21 -1.14 -6.34	-,34 -,25 -,29	-11.39 5.89 -1.14 -6.64	-11.39 4.47 -7.28 -14.21	.38 95 57	-11.39 4.84 -8.23 -14.78

Figure 20 (Continued)

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Figure 21. Binding energy decomposition and description of electron distribution for HF, BAO approximation

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BINDING ENERGY PARTITIONING

	1	Prom	otion	Qua	asi- sical	Sha Pene	ring tration	Shari Interíe	ing rence	
		PRH	PRC	QC N	QCT	SPN	SPT	SIN	SIT	TOTAL
F	KIN NA OEI TOT	2.79 .17 56 2.41			16.38 -14.98 .62 2.02	ό.ύ4 ό.υ4	-2.75 -2.75			12, 17 -14, 81 3, 95 8, 32
н	KIN SNA OEI TOT	-	-		- 2.98 2.24 .41 34	3. 33 3. 33	2.01 2.01			- 2.98 2.24 5.75 5.00
BOND	KIN SNA OEI TOT			-1.83 -1.83	. 48 70 .28	-0.00 -0.00	04 04	-10.75 5.71 - 1.04 - 6.07	41 41	-10.75 4.44 - 8.37 -14.67
HF	KIN SNA OEI TOT	2,73 ,17 - ,56 2,41	-	-1.83	13.40 -11.76 .33 1.26	3. 17 3. 17	-1.38 -1.38	-10.75 5.71 - 1.04 - 6.07	41 41	5.44 - 8.12 1.32 - 1.36

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DENSITY PARTITIONING FOR VALENCE AO'S

	Vale	nce Atomic Or	rbitals	Overlap		ons				
	ĸ	5	J	S(F.H)	ę	v	q	ç ^N	ΥP	^p (F, H)
1	. 991851	126137	. 017 20	. 000000	2.000	. 000	2.000	2. 000	. 000	. 000
٤	. 124571	. 730657	344030	. 346572	2.233	171	2.062	1. 264	.078	492
ь	. 026719	. 343404	. 238780	. 454322	. 850	. 307	1.157	1.010	. 141	. 675
#		(unhybridize)	d)	0	2.000	. 000	2.000	2.000	. 000	. 000
h		(unhybridize)	d)	-	. 645	. 136	. 781	1.000	219	-

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

			·····	F		н
		i	g	ь	17	h
	PR	2.000	. 000	. 000	. 000	-
	SPN	. 000	. 000	. 000	.000	. 000
1	SPT	. 000	. 000	. 000	. 000	. 000
	TOT	2.000	. 000	. 000	.000	.000
	PR	. 000	2.004	020	. 000	-
•	SPN	. 000	. 000	. 001	. 000	001
x	SPT	. 000	. 158	. 006	. 000	085
	TOT	. 000	2.162	013	.000	087
	PR	. 000	020	1.037	. 000	-
	SPN	. 000	. 001	395	. 000	. 394
b	SPT	. 000	.006	. 0 32	. 000	. 102
	тот	. 000	013	. 674	. 000	. 496
	PR	. 000	.000	. 000	2.000	-
	SPN	. 000	. 000	. 000	. 000	.000
17	SPT	. 000	. 000	. 000	. 000	.000
	тот	. 000	. 000	. 000	2.000	. 000
	PR	-	-	-	· _	1.000
	SPN	. 000	001	. 394	. 000	392
h	SPT	. 000	085	. 102	.000	236
	TOT	. 000	087	. 496	. 000	. 372

							Quasi- Sharing classical Penetration			ring	SP+QC
Grbital	×	1	. or 1	TT TT	₩	Total	Orbital	QCT	SPN	SPT	Total
F k OEI TOT	97 2.02 20 .85	3.84 -3.17 -4.22 -4.35	.00 .00 4.50 4.50	.20 .20 23 03	.02 .02 03 23	93 1.97 43 .61	F 1 OEI TOT	00. 99. 59. 59.	62 20	.ee .ec	. 00 . 02 00 00
KIN NA P = OEI TOT	3.84 -3.17 -4.22 -4.35	-13.91 45.40 -3.62 27.86	.00 .00 .15	.00 .00 -3.33 -3.33	.22 .20 - <i>3.33</i> -3.33	-7.87 39.11 -25.67 5.57	P P SNA P P OEI TOT	4.24 -4.84 .21 38	.62 .62	-2.21 -2.21	4.24 -4.84 -1.40 -1.99
F or OEI TOT	.00 .00 4.50 4.50	.00 .00 .15 .15	11.59 -40.91 3.40 -25.92	.20 .00 3.17 3.17	.00 .00 3.17 3.17	11.59 -40.91 25.75 -3.57	F b OEI TOT	12.13 -10.14 .41 2.41	6.04 6.04	54 54	12.13 -10.14 5.91 7.91
KIN NA P TT OEI TOT	.00 03 03	.00 .00 -3.33 -3.33	.02 .00 3.17 3.17	.00 10 10	. 80 . 80 . 80 . 82	.00 .00 10 10	KIN P T OEI TOT	.08 .00 .00	00 00	.00 .00	. ee . ee ea ea
P TT OEI TOT	.00 .00 03 03	.00 .00 -3.33 -3.33	.00 .00 3.17 3.17	. 88 . 88 . 88 . 88	.00 .00 10 10	.00 .00 12 12	P T SNA QEI TOT	.00 .02 .00 .80	00 00	.00 .00	99 90 - 92 - 02
Total P KIN NA OEI TOT						2.79 .17 56 2.41	Total F KIN SNA OEI TOT	16.38 -14.98 .62 2.82	6.64 6.64	-2.75 -2.75	16.38 -14.98 4.51 5.91
							Total H KIN SNA OEI TOT	-2.98 2.24 .41 34	3.33	2.81	-2.98 2.24 5.75 5.20

INTEA-ATOMIC ENERGY CONTRIBUTIONS OF HP, BAO

INTERATORIC	ENERGY	CUNTRIBUTIONS	0P	HP.	BAO

Orbit	al	Qua	siclassical Electro	static			I	Totals					
Pal P R	r	QCN	QCN QCT		SPN	SPN SPT		SIN	- Interference	SI(N+T)	I (N+T) N	т	N+T
1 h	KIN SNA OEI TOT	05	.00 .01 .00 .01	04 .00 04	.00	. ee . ee	.83 .99	.10 .00 29 .01	.20 .80	. ee . 19 89 . 01	.22 .25 29 84	.01 .00 .01	.00 .06 09 02
L h	KIN SNA OEI TOT	1.12	86 1.89 28 .82	2.15 28 1.94	• 18 • 18	1. 01 1. 01	1.19 1.19	10.30 -1.3476 11 8.10	1503 18	10.30 -2.28 11 7.92	10.30 98 .27 9.39	.85 ,81 1.66	12.32 13 .88 11.05
bh	KIN SNA OEI TOT	-3.27 -3.27	6439 50 -1.52	-4.30 50 -4.79	-6.17 -6.17	-1.65 -1.65	-7.82 -7.82	-21.05 9.21 -2.39 .15 -14.06	.1e33 23	-21.05 6.58 .15 -14.31	-21.85 3.55 -6.82 -23.52	-1.26 -2.15 -3.41	-21.25 2.29 -8.17 -26.93
Ψh	KIN SNA OEI TOT	•18 •18	.80 .43 .00 .48	.67 .20 .67	. 20 . 22	. 82 . 82	. ee	.eə .45 .ee 58 65	.00 .00. DD	. 22 . 45 52 25	.00 .63 50 .13	.48 .00 .45	.00 1.11 50 .62
ሞ ከ	KIN SNA OEI TOT	•18 •18	.80 .48 .82 .48	.67 .00 .67	.00 .02	. ee . ce	. ee . ee	.0C .45 .00 50 05	99. 99. 96 .	.00 .45 50 05	.02 .63 58 .13	.48 .02 .48	.02 1.11 50 .62
Bond Total	KIN SNA OEI TOT	-1.83	79 1.68 70 .28	85 70 -1.55	-6.00	64	-6.63 -6.63	-10.75 8.86 -3.15 -1.04 -6.07	2536 41	-12.75 5.32 -1.24 -6.49	-12.75 3.88 -7.23 -13.90	.57 -1.34 77	-10.75 4.44 -8.37 -14.67

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Figure 21 (Continued)

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Figure 22. Binding energy decomposition and description of electron distribution for HF, BMAO approximation

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BINDING ENERGY PARTITIONING

		Promotion		Qua	asi- sical	Sha Pene	ring tration	Shari Interfe	ing			
		PRH	PRC	QCN	QCT	SPN	SPT	SIN	SIT	TOTAL		
F	KIN NA OEI TOT	2.33 .59 2.67 5.59	- 5,50 10.75 - 5.08 .17		19.03 -17.14 .71 2.60	- 3, 92 - 3, 92	5.29 5.29			15.87 - 5.80 33 9.72		
н	KIN SNA OEI TOT	. 00 . 00 . 00 . 00	9,97 - 8,61 .00 1,36		- 5.41 3.08 .57 - 1.74	-1,35 -1,35	8.69 8.69			4.56 - 5.53 7.94 6.97		
BOND	KIN SNA OEI TOT			-2.07 -2.07	1.39 85 .54	3. 48 3. 48	-11.07 -11.07	-17.36 8.37 77 - 9.76	36 36	-17.36 7.33 - 9.21 -19.25		
HF	KIN SNA OEI TOT	2.33 .53 2.67 5.53	4.47 2.14 - 5.08 1.53	-2.07 -2.07	13.62 -12.67 .45 1.40	-1.79 -1.79	2.91 2.91	-17.36 8.37 77 - 9.76	36 36	3,06 - 4,00 - 1,61 - 2,55		

DENSITY PARTITIONING FOR VALENCE AO'S

	Valer	nce Atomic Or	bitals	Overlap								
	k	s	σ	S _(F,H)	p	v	q	q ^N	q ^T	^p (F, H)		
i	. 991032	1270 ??	. 022022	. 000000	2.000	. 000	2.000	2.000	. 000	. 000		
l	.127041	. 429591	346008	. 299000	2.166	120	2.040	1.982	. 004	401		
ь	. J22 - 48	. 345987	, 937959	. 467125	. 873	. 310	1.184	1.018	. 166	. 665		
11		(unhybridize)	d)	0	2.000	. 000	2.000	2.000	. 000	. 000		
h		(unhybridize)	d)	-	. 580	. 191	.771	1.000	230	-		

EXCHANGE CONTRIBUTIONS TO PAIRPOPULATIONS

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			F								
		i	L	ь	n	h					
•	PR	2.000	. 000	. 000	. 000	-					
	SPN	. 000	. 000	.000	. 000	. 000					
1	SPT	. 000	. 000	. 300	.000	. 000					
	тот	2.000	. 000	.000	. 000	. 000					
	PR	. 000	1.502	. 479	. 000	-					
	SPN	. 000	. 308	501	.000	. 193					
£	SP'I	. 000	. 306	.013	. 000	254					
	тот	. 000	2.117	009	.000	062					
	PR	.000	. 479	. 539	. 000	-					
	SPN	. 000	501	. 814	. 000	313					
Þ	SPT	. 000	. 013	649	. 000	. 802					
	тот	.000	009	. 704	. 000	. 489					
	PR	. 000	. 000	. 000	2.000	-					
	SPN	. 000	. 000	. 000	.000	. 000					
11	SPT	. 000	.000	. 000	. 000	. 000					
	TOT	. 000	. 000	.000	2.000	.000					
				4							
	PP	_	-	-	-	1,000					
	SPN	000	193	- 313	. 000	. 120					
h	SPT	000	- 254	802	. 000	777					
	TOT	.000	- 067	. 489	. 000	. 343					
	• • • •		,	/							

Figure 22 (Continued)

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INTEA-ATOMIC ENERGY CONTRIBUTIONS OF HP, BMAD

					- Prosot	reaction [1	Quas1-	Sharing		SP+QC			
Orbital	PBH	PBC	PRH	PBC	PRH	PBC	PRH	T PBC	PRH	FRC PRC	PRH	PRC	PB(B+C)	Orbital	Classical QCT	SPN SPN	SPT	Total
KIN F k OBI TOT	83 1.72 13 77	1.51 -1.57 .89 .03	2.78 -2.89 -4.28 -4.39	81 .62 21 21	.88 .58 4.52 4.52	. 80 . 80 1.75 1.75	.00 .90 03 03	.00 .00 -1.38 -1,38	.00 .00 03 03	.00 .00 -1.38 -1.38	79 1.69 32 .57	1.58 -1.68 .25 .23	.78 .89 07 .50	P 1 OEI TOT	. 80 . 80 . 80 . 80	00 ~.00	.90 .80	- 68 - 68 - 68
F = OEI TOT	2.75 -2.89 -4.28 -4.39	81 .02 21 21	-14.26 46.53 .46 32.73	-1.97 2.78 03 .78	.90 -2.80 -2.99	.00 .00 .42 .42	.00 -3,33 -3,33	.00 .00 52 52	.00 .00 -3.33 -3.33	.08 .98 52 52	-8.74 40.75 -14.56 17.49	-2.06 2.84 92 14	-10.80 43.62 -15.48 17.34	P & SWA OEI TOT	3.45 -3.99 .12 36	2.18 2.18	-4.32 -4.32	3.45 -3.99 -1.96 -2.49
F or OEI TOT	.88 .88 4.52 4.52	.00 .00 1.75 1.75	.00 .00 -2.80 -2.80	-00 -00 -42 -42	11.87 -41.88 7.97 -22.05	9.62 -16.59 .38 -6.67	.00 .00 3.24 3.24	.00 .00 .31 .31	.00 .20 3.24 3.24	.00 .31 .31	11.97 -41.98 17.60 -12.41	9.62 -16.59 7.81 .84	21.49 -58.47 24.62 -12.37	F b OEI TOT	15.57 -13.15 .54 2.95	-6.11 -6.11	9.61 9.61	15.57 -13.15 4.84 6.45
KIN PT OEI TOT	.00 .00 03 03	.00 .00 -1.38 -1.38	.00 .00 -3.33 -3.33	.80 .00 52 52	.00 .00 3.24 3.24	.00 .00 .31 .31	.00 .00 03 03	-7.32 13.85 50 5.23	.00 .00 .00	.00 .99 -1.90 -1.90	.90 .00 93 93	-7,32 13,05 -5,71 ,02	-7.32 13.85 -5.74 01	F TT SHA DEI TOT	.00 .00 .00	08 08	.00 .02	.08 .09 09 09
P TT NA OEI TOT	.00 .00 03 03	.00 .00 -1.38 -1.38	.00 .00 -3.33 -3.33	.00 .00 52 52	.80 .89 3.24 3.24	.00 .00 .31 .31	00. 88 00 00	.88 .98 -1.98 -1.00	.00 .00 03 03	-7.32 13.85 -50 5.23	.90 .00 03 03	-7.32 13.85 -5.71 .82	-7.32 13.05 -5.74 01	P # SNA OEI TOT	.80 .80 .90 .88	00 00	.88 .88	88 99 - 99 - 99
Total P KIN NA OBI TOT											2.55 .59 2.67 5.59	-5.50 10.75 -5.69 .17	-5.16 11.34 -2.41 5.76	Total P KIN SNA OEI TOT	19.03 -17.14 .71 2.60	-3.92 -3.92	5.29 5.29	19.03 -17.14 2.08 3.96
Total H KIN NA OEI TOT											.00 .00 .00 .00 .00	9,97 -8,61 .00 1.36	9.97 -8.61 .28 1.36	Total H KIN SNA OEI TOT	-5.41 3.08 .59 -1.74	-1.35	8.69 8.69	-5.41 3.08 7.94 5.61

INTERATORIC ENERGY CONTRIBUTIONS OF HP, BRAD

Orbita	11	Quas	iclassical Electro	static	,		Sh	aring		·	Totale			
Pair F H		QCH	QCT	QC(N+T)	SPN	SPT	SP(N+T)	SIN	SIT	SI(N+T)	и	T	N+T	
1 h	KIN SHA OEI TOT	84 84	.00 .01 .88 .81	83 83	.08 .00	. 20 . 28	. 80 . 80	.08 .0 06 .02	89. 89. 89. 99.	.00 .08 06 .02	.00 .04 86 02	.01 .00 .61	.08 .85 06 21	
£ h	KIN SNA OEI TOT	.37 .37	02 1.29 18 1.89	1.65 18 1.47	-2.22 -2.22	3.13	.91 .91	8.70 -1.3580 11 6.38	50905 14	8.70 -2.36 11 6.23	8.70 -1.84 -2.33 4.54	1.13 2.95 4.08	8.70 71 .63 8.61	
bh	KIN SNA OEI TOT	-2.30 -2.30	4660 67 -1.73	-3.36 67 -4.04	5.69	-14.28 -14,28	-8.51 -8.51	-26.86 11.3284 -15.48	• • • • • • • • • • • • • • • • • • •	-26.P6 10.24 .12 -15.70	-26.06 8.16 5.91 -12.09	-1.28 -14.97 -16.15	-26.06 6.88 -9.06 -28.24	
η p	KIN SNA OEI TOT	85 85	.00 .58 .00 .58	.53 .68 .53	. 00 . 68	. 20 . 20	. 60 . 64	.82 .81 36 33	89. 89. 89.	.02 36 33	03 36 39	.58 .00 .58	.80 .55 36 .20	
πħ	KIN SNA OEI TOT	85 05	.C0 .58 .80 .58	.53 .80 .53	.00 .00	. 00 . 00	. 00 . 00	.62 .08 .6 35 33	99.68 .69	.88 .82 36 33	.08 83 36 39	. 53 . <i>88</i> . 58	.00 .55 36 .20	
Bond Total	KIN SNA OEI TOT	-2.07	~.49 1.86 85 .54	68 85 -1.54	3.48 3.48	-11.67 -11.67	-7.59	-17.36 18.89 -1.7: -,77 -9.76	 20234 36	-17,36 8.01 77 -18,12	-17.36 6.30 2.71 -8.35	1.03 -11.92 -10.89	-17.36 7.33 -9.21 -19.25	

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